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AUG 77 E L BELL, D K COHOON, J W PENN

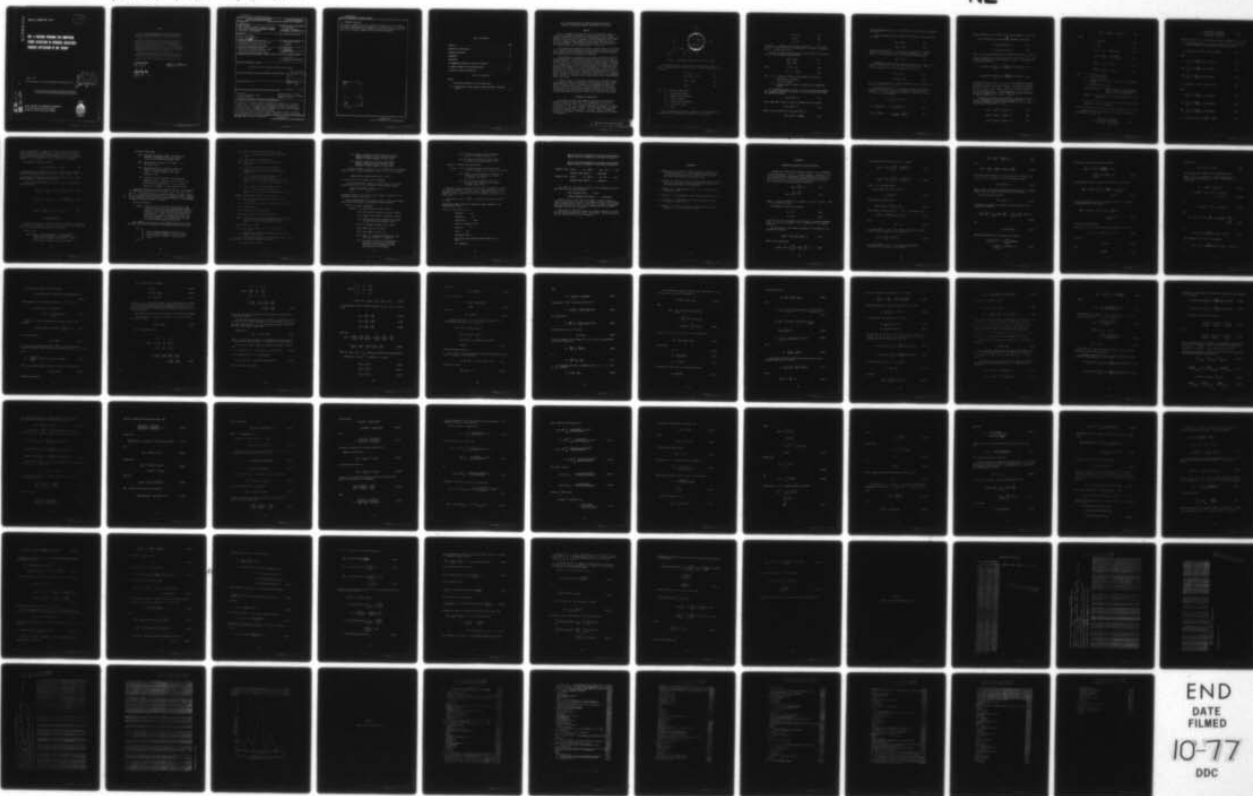
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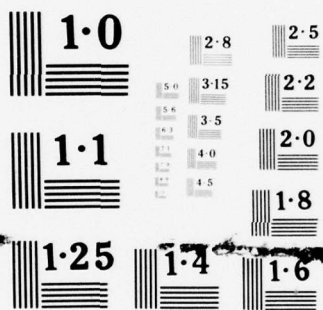
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**MIE: A FORTRAN PROGRAM FOR COMPUTING
POWER DEPOSITION IN SPHERICAL DIELECTRICS
THROUGH APPLICATION OF MIE THEORY**

August 1977

Interim Report for Period September 1976-January 1977

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This report has been reviewed by the Information Office (OI) and is releasable to the National Technical Information Service (NTIS). At NTIS, it will be available to the general public, including foreign nations.

This technical report has been reviewed and is approved for publication.

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20. ABSTRACT (Continued)

structure and sequence of control parameter and data cards, output format, and function subprograms and subroutines are covered in detail. An extensive discussion of the Mie solution, sample problems with associated computer results, and a listing of program MIE are included in appendixes.

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TABLE OF CONTENTS

	Page
PURPOSE	3
MATHEMATICAL DESCRIPTION	3
PROGRAM DESCRIPTION	12
REFERENCES	18
APPENDIXES	
A--MATHEMATICAL ANALYSIS OF THE MIE SOLUTION	19
B--SAMPLE PROBLEMS WITH COMPUTER RESULTS	62
C--SOURCE LISTING OF PROGRAM MIE	71

List of Illustrations

Figure

1 Directional Approach of the Incident Wave	4
B-1 Distribution of Power Density Inside the Sphere Along the z Axis	70

MIE: A FORTRAN PROGRAM FOR COMPUTING POWER DEPOSITION IN SPHERICAL DIELECTRICS THROUGH APPLICATION OF MIE THEORY

PURPOSE

MIE is a FORTRAN IV program written for the IBM 360/65 system, to compute the absorbed power density at internal points, the average absorbed power density, and the total absorbed power inside a homogeneous spherical dielectric that is immersed in an electromagnetic field. Results are obtained through applying the Mie theory and using double-precision arithmetic. Here double-precision numbers have approximately 16.8 decimal digits and an exponent range of -78 to +75.

This program was originally produced to provide special test results for a computer program designed to solve the problem of ascertaining the power deposition inside arbitrarily shaped, finitely conducting biological bodies exposed to electromagnetic radiation.

The knowledge to be gleaned from this effort is directly related to the research effort of the Radiation Sciences Division at the School of Aerospace Medicine. Briefly, here studies are being currently conducted (1) to determine the radiofrequency radiation-induced effects in biological specimens, (2) to seek out possible hazards to personnel in a radiofrequency environment, (3) to accurately measure and to determine the distribution of energy in the whole biological body or just in a particular organ, (4) to find ways to reduce any potentially adverse action between RF emitters and cardiac pacemaker and prosthetics, (5) to extrapolate response to radiation from the test animal to man in a meaningful manner, and (6) to contribute in the design of realistic safety standards with a solid basis.

To benefit users of this report, program MIE is described in sufficient depth to facilitate implementation on any modern computer and job setups. Detailed coverage is provided of the mathematical theory and formulas, structure and sequence of control parameter and data cards, output format, and function subprograms and subroutines. Included in the appendixes are an extensive discussion of the Mie solution, sample problems with associated computer results, and a listing of the program.

MATHEMATICAL DESCRIPTION

We consider the plane electromagnetic wave that is irradiating a homogeneous spherical dielectric to be propagating in the positive z -direction, and the electric field E to be linearly polarized in the x -direction (cf. Fig. 1). A system of Cartesian coordinates with origin at the center of the sphere is used. Also, the medium which surrounds the sphere is taken as free space (or vacuum). Thus our embedding medium is a nonconductor, and both the surrounding medium and the sphere are nonmagnetic.

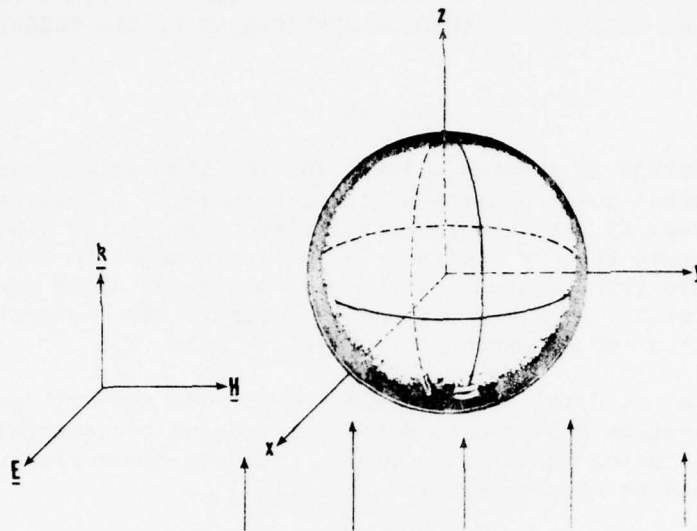


Figure 1. Directional approach of the incident wave.

The Gaussian system of units (nonrationalized c.g.s. units) is used. In this system the macroscopic form of Maxwell's equations may be written as

$$\nabla \cdot \underline{D} = 4\pi\rho, \quad (1)$$

$$\nabla \times \underline{H} = \frac{4\pi}{c} \underline{J} + \frac{1}{c} \frac{\partial \underline{D}}{\partial t}, \quad (2)$$

$$\nabla \times \underline{E} = -\frac{1}{c} \frac{\partial \underline{H}}{\partial t}, \quad (3)$$

$$\nabla \cdot \underline{H} = 0, \quad (4)$$

with \underline{D} = electric displacement,
 ρ = free-charge density,
 \underline{H} = magnetic-field intensity,
 \underline{J} = current density,
 \underline{E} = electric-field intensity,
 c = velocity of light,
 t = time.

For our homogeneous, isotropic, permeable, conducting dielectric, and linear case, the constitutive (material) relations are

$$\underline{D} = \epsilon \underline{E} , \quad (5)$$

$$\underline{B} = \underline{H} , \quad (6)$$

$$\underline{J} = \sigma \underline{E} . \quad (7)$$

The symbol ϵ represents the dielectric constant, and σ the conductivity. In these relations and in equations 1-4, the magnetic permeability, μ , was set equal to one.

Since we are considering time-harmonic fields, Maxwell's equations 2 and 3 take the simpler time-free form

$$\nabla \times \underline{H} = im^2 k \underline{E} , \quad (8)$$

$$\nabla \times \underline{E} = -ik \underline{H} , \quad (9)$$

where

$$k = \frac{2\pi}{\lambda} = \frac{\omega}{c} , \quad (10)$$

$$m^2 = \epsilon - i \frac{4\pi\sigma}{\omega} , \quad (11)$$

with k = propagation constant or wave number in free space,
 λ = wavelength in free space,
 ω = circular frequency,
 m = complex refractive index, a constant for a homogeneous medium.

In a homogeneous medium, relative to a Cartesian coordinate system, each rectangular component of vectors \underline{E} and \underline{H} satisfies the scalar-wave Helmholtz equation

$$\nabla^2 \psi + m^2 k^2 \psi = 0. \quad (12)$$

At the same time, vectors \underline{E} and \underline{H} satisfy the vector equation

$$\nabla^2 \underline{A} + m^2 k^2 \underline{A} = 0. \quad (13)$$

Recall that the formula for evaluating $\nabla^2 \underline{A}$ is

$$\nabla^2 \underline{A} = \nabla(\nabla \cdot \underline{A}) - \nabla \times (\nabla \times \underline{A}). \quad (14)$$

(Details showing that vector \underline{E} satisfies eq. 13 can be found in Appendix A, pp. 19-21.)

Now two independent vector solutions of equation 13 can be expressed as vectors

$$\underline{M}_{\psi} = \nabla \times (\underline{r}\psi), \quad (15)$$

$$mk\underline{N}_{\psi} = \nabla \times \underline{M}_{\psi}, \quad (16)$$

where the scalar function ψ is a solution of equation 12. (Details pertinent to eqs. 15 and 16 can be found in Appendix A, pp. 27-30.) Since \underline{r} is a constant vector, the following relationship holds:

$$mk\underline{M}_{\psi} = \nabla \times \underline{N}_{\psi} \quad (17)$$

Consequently, if we choose two independent solutions, u and v , of equation 12 and construct the field vectors $\underline{M}_u, \underline{M}_v, \underline{N}_u, \underline{N}_v$, we will find that equations 8 and 9 are satisfied by the vectors

$$\underline{E} = \underline{M}_v + i\underline{N}_u, \quad (18)$$

$$\underline{H} = m(-\underline{M}_u + i\underline{N}_v). \quad (19)$$

Here i is the complex unit. (Details pertinent to eqs. 18 and 19 can be found in Appendix A, pp. 30-31.)

Convenience dictates that our problem be embedded in a spherical system of coordinates r (radial), θ (scattering or colatitude), and ϕ (azimuthal). In this system, the components of vectors \underline{M}_{ψ} and \underline{N}_{ψ} ($\psi = u$ or v) can be expressed as

$$\underline{M}_r = 0, \quad \underline{N}_r = \frac{1}{mk} \frac{\partial^2}{\partial r^2} (r\psi) + mkr\psi, \quad (20)$$

$$\underline{M}_{\theta} = \frac{1}{r \sin \theta} \frac{\partial}{\partial \phi} (r\psi), \quad \underline{N}_{\theta} = \frac{1}{mkr} \frac{\partial^2}{\partial r \partial \theta} (r\psi), \quad (21)$$

$$\underline{M}_{\phi} = -\frac{1}{r} \frac{\partial}{\partial \theta} (r\psi), \quad \underline{N}_{\phi} = \frac{1}{mkr \sin \theta} \frac{\partial^2}{\partial r \partial \phi} (r\psi). \quad (22)$$

(Details pertinent to eqs. 20-22 can be found in Appendix A, pp. 32-34.)

For our incident plane-wave radiation expressed in the form

$$\underline{E} = \underline{a}_x E_0 \exp[-i(kz - \omega t)], \quad (23)$$

$$\underline{H} = \underline{a}_y E_0 \exp[-i(kz - \omega t)], \quad (24)$$

where E_0 is the wave amplitude, and polarization vectors \underline{a}_x and \underline{a}_y are unit vectors directed along the x,y-axes, van de Hulst (5, p. 122) chooses the independent scalar solutions, u and v , of equation 12 for the induced wave as

$$u = E_0 \exp(i\omega t) \cos\phi \sum_{n=1}^{\infty} m c_n (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos\theta) j_n(mkr), \quad (25)$$

$$v = E_0 \exp(i\omega t) \sin\phi \sum_{n=1}^{\infty} m d_n (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos\theta) j_n(mkr), \quad (26)$$

The expansion coefficients c_n and d_n are determined by using the appropriate boundary conditions; $P_n^1(\cos\theta)$, the associated Legendre function of order 1; and $j_n(mkr)$, the spherical Bessel function of the first kind and order n . (Details pertinent to eqs. 25 and 26 and the corresponding forms for the incident and scattered waves can be found in Appendix A, pp. 34, 36, 37, 49-61.)

The boundary conditions on the tangential components of vectors \underline{E} and \underline{H} , when applied to the components in equations 20-22, yield a system of algebraic equations in the unknowns a_n , b_n , c_n , and d_n (van de Hulst's notation [5, p. 123]):

$$\psi_n(x) - a_n \zeta_n(x) - m c_n \psi_n(y) = 0, \quad (27)$$

$$\psi'_n(x) - a_n \zeta'_n(x) - c_n \psi'_n(y) = 0, \quad (28)$$

$$\psi_n(x) - b_n \zeta_n(x) - d_n \psi_n(y) = 0, \quad (29)$$

$$\psi'_n(x) - b_n \zeta'_n(x) - m d_n \psi'_n(y) = 0, \quad (30)$$

where

$$x = \frac{2\pi a}{\lambda} = ka, \quad (31)$$

$$y = mka, \quad (32)$$

$$\psi_n(z) = z j_n(z) = \left(\frac{\pi z}{2}\right)^{1/2} J_{n+1/2}(z), \quad (33)$$

$$\zeta_n(z) = z h_n^{(2)}(z) = \left(\frac{\pi z}{2}\right)^{1/2} H_{n+1/2}^{(2)}(z), \quad (34)$$

$$h_n^{(2)}(z) = j_n(z) - i n_n(z), \quad (35)$$

$$H_n^{(2)}(z) = J_{n+1/2}(z) - i N_{n+1/2}(z). \quad (36)$$

Here x = MIE-size parameter,
 a = radius of the sphere,
 m = complex refractive index, a constant for a homogeneous medium,
 λ = wavelength in free space,
 $J_{n+1/2}(z)$ and $N_{n+1/2}(z)$ = Bessel functions of half-odd-integer order, of the first and second kind (Neumann function) respectively,
 $H_{n+1/2}^{(2)}(z)$ = Hankel function of half-odd-integer order, of the second kind.

The prime on $\psi_n(z)$ or $J_n(z)$ denotes first derivative with respect to the function argument.

(Details pertinent to eqs. 27-36 can be found in Appendix A, pp. 35-45.)

Elementary algebraic manipulations of equations 27-30 result in the solution of c_n and d_n :

$$c_n = \frac{\psi'_n(x) \zeta_n(x) - \psi_n(x) \zeta'_n(x)}{\psi'_n(y) \zeta_n(x) - m \psi_n(y) \zeta'_n(x)}, \quad (37)$$

$$d_n = \frac{\psi'_n(x)\zeta_n(x) - \psi_n(x)\zeta'_n(x)}{m\psi'_n(y)\zeta_n(x) - \psi_n(y)\zeta'_n(x)}. \quad (38)$$

(Details pertinent to eqs. 37 and 38 can be found in Appendix A, pp. 42-49.)

Utilizing equations 18, 20, 21, 22, 25, and 26, we can derive the r , θ , and ϕ components for the induced electric-field vector:

$$E_r = iE_0 \exp(i\omega t) \cos\phi (m^2 krU + 2mUR + m^2 krURR), \quad (39)$$

where

$$U = \sum_{n=1}^{\infty} c_n (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos\theta) j_n(mkr), \quad (40)$$

$$UR = \sum_{n=1}^{\infty} c_n (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos\theta) j'_n(mkr), \quad (41)$$

$$URR = \sum_{n=1}^{\infty} c_n (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos\theta) j''_n(mkr). \quad (42)$$

$$E_\theta = E_0 \exp(i\omega t) \cos\phi [mU + i(\frac{1}{kr}UR + mURR)] \quad (43)$$

where

$$U = \sum_{n=1}^{\infty} d_n (-i)^n \frac{2n+1}{n(n+1)} \pi_n(\cos\theta) j_n(mkr), \quad (44)$$

$$UR = \sum_{n=1}^{\infty} c_n (-i)^n \frac{2n+1}{n(n+1)} \tau_n(\cos\theta) j_n(mkr), \quad (45)$$

$$URR = \sum_{n=1}^{\infty} c_n (-i)^n \frac{2n+1}{n(n+1)} \tau_n(\cos\theta) j'_n(mkr). \quad (46)$$

$$E_\phi = -E_0 \exp(i\omega t) \sin\phi [mU + i(\frac{1}{kr}UR + mURR)], \quad (47)$$

where

$$U = \sum_{n=1}^{\infty} d_n (-i)^n \frac{2n+1}{n(n+1)} \tau_n(\cos\theta) j_n(mkr), \quad (48)$$

$$UR = \sum_{n=1}^{\infty} c_n (-i)^n \frac{2n+1}{n(n+1)} \pi_n(\cos\theta) j_n(mkr), \quad (49)$$

$$URR = \sum_{n=1}^{\infty} c_n (-i)^n \frac{2n+1}{n(n+1)} \pi_n(\cos\theta) j'_n(mkr). \quad (50)$$

Here a single or double prime denotes a first or second derivative with respect to mkr and

$$\pi_n(\cos\theta) = \frac{1}{\sin\theta} P_n^1(\cos\theta), \quad (51)$$

$$\tau_n(\cos\theta) = \frac{d}{d\theta} P_n^1(\cos\theta). \quad (52)$$

For a sphere of radius a (cm), the absorbed power density at an internal point, the average absorbed power density, and the total absorbed power are computed by means of the following formulas (with σ and \underline{E} in electrostatic units of the nonrationalized c.g.s. system):

$$P_d(r, \theta, \phi) = 0.05\sigma \underline{E} \cdot \underline{E}^* \quad \text{watts/m}^3, \quad (53)$$

$$P_{\text{tot}} = 10^{-6} \int_0^\pi \int_0^a \int_0^{2\pi} P_d r^2 \sin\theta d\phi dr d\theta \quad \text{watts}, \quad (54)$$

$$P_{\text{avg}} = 10^6 P_{\text{tot}} / [(4/3)\pi a^3] \quad \text{watts/m}^3. \quad (55)$$

Here, $*$ denotes the complex conjugate. In program MIE a numerical integration product rule, based on three m -point Gauss-Legendre quadrature formulas, is used to evaluate the triple integral.

To complete our discussion, it seems appropriate to consider the formulas used in generating the values of certain functions. The formulas

$$P_{n+1}^1(\cos\theta) = \frac{2n+1}{n} \cos\theta P_n^1(\cos\theta) - \frac{n+1}{n} P_{n-1}^1(\cos\theta), \quad (56)$$

$$\sin\theta \frac{d}{d\theta} P_n^1(\cos\theta) = n \cos\theta P_n^1(\cos\theta) - (n+1)P_{n-1}^1(\cos\theta), \quad (57)$$

together with

$$P_1^1(\cos\theta) = \sin\theta, \quad (58)$$

$$P_2^1(\cos\theta) = 3 \sin\theta \cos\theta \quad (59)$$

are used to generate function and derivative values of the associated Legendre functions.

Special limit values are also obtained by

$$\lim_{\theta \rightarrow 0} \frac{P_n^1(\cos\theta)}{\sin\theta} = \frac{n(n+1)}{2}, \quad (60)$$

$$\lim_{\theta \rightarrow \pi} \frac{P_n^1(\cos\theta)}{\sin\theta} = \frac{(-1)^{n+1} n(n+1)}{2}. \quad (61)$$

The forward recurrence relation

$$\eta_{n+1}(z) + \eta_{n-1}(z) = \frac{2n+1}{z} \eta_n(z) \quad (62)$$

is used together with the relations

$$\eta_0(z) = -\frac{\cos z}{z}, \quad (63)$$

$$\eta_1(z) = -\frac{\cos z}{z^2} - \frac{\sin z}{z} \quad (64)$$

to generate values of the Neumann spherical Bessel functions. The generating process is terminated at order N when the following termination criterion

$$|\eta_n(z)| \geq \text{STOPR} \quad (65)$$

is met. Here STOPR is a number, say 1.0D25. The user's needs will determine whether or not STOPR should retain its presently assigned value. Our own demands were satisfactorily met for both real ($2\pi r/f$) and complex ($2\pi mr/f$) arguments of the Neumann functions for parameter ranges: $7 \leq |m| \leq 100$, $5 \leq r \leq 25$ cm, and $20 \leq f \leq 1000$ MHz.

The backward recurrence relation

$$j_{n-1}(z) = \frac{2n+1}{z} j_n(z) - j_{n+1}(z) \quad (66)$$

in combination with an appropriate starting value is used to generate values of the spherical Bessel functions of the first kind, $j_n(z)$.

This technique of using the backward relation in place of the forward relation helps to avoid stability problems.

Values of the derivatives of $j_n(z)$ and $\eta_n(z)$ are obtained by using the formulas:

$$\frac{d}{dz} j_n(z) = \frac{1}{2n+1} [n j_{n-1}(z) - (n+1) j_{n+1}(z)], \quad (67)$$

$$\frac{d^2}{dz^2} j_n(z) = -\frac{2}{z} \frac{d}{dz} j_n(z) - \left[1 - \frac{n(n+1)}{z^2}\right] j_n(z), \quad (68)$$

$$\frac{d}{dz} \eta_n(z) = \frac{1}{2n+1} [n \eta_{n-1}(z) - (n+1) \eta_{n+1}(z)]. \quad (69)$$

PROGRAM DESCRIPTION

Program MIE consists of a driver routine, four subroutine subprograms, and one function subprogram. A list of the driver routine and subprograms, along with a brief explanation of their use, follows.

Driver routine:

MAIN - used to input/output data, to compute spherical Bessel function values for a real argument, to compute series expansion coefficients c_n and d_n , and to direct the course of the calculations.

Subroutine subprograms:

EVEC - generates the complex radial, colatitude, and azimuthal components of the electric-field vector, \underline{E} , and the scalar product $\underline{E} \cdot \underline{E}^*$.

TERM - generates the n th power of $-i$, where i is the complex unit.

BESS - generates an array of values of each of the spherical Bessel functions $j_n(mx)$ and $\eta_n(mx)$ for complex argument mx . Array dimension = 100.

PL - generates an array of values of the associated Legendre function, $P_n^1(\cos\theta)$, and an array for its first derivative, $\frac{d}{d\theta} P_n^1(\cos\theta)$, for n varying from one to a maximum N . Array dimension = 100.

Subroutines BESS and PL, plus the algorithm (in-line coded in MAIN) for computing the spherical Bessel functions $j_n(x)$ and $\eta_n(x)$ for real x , are modified versions of subroutines found in program SUP[1]-- one of two allied programs developed by D. S. Drumheller and D. E. Setzer (1) as a research tool in the study of problems related to electromagnetic transmission through atmospheric aerosols.

Function subprogram:

GAUSS3 - a product rule for the numerical evaluation of the triple integral for the total absorbed power within the sphere. A basic m -point Gauss-Legendre quadrature formula is used. User's option is available for the selection of m , number of weighting coefficients, and associated points for Gaussian quadrature, from the set of integers $\{2,3,4,5,6,8,10,12,14\}$ for each of the three Gauss quadrature rules used.

Blank COMMON is used by driver routine (MAIN) and subroutine EVEC. The list of the arrays, variables, and constants stored in this area is

CN -
 }
DN - array of complex expansion coefficients for
 coefficients for components of electric-field
 vector \underline{E} inside the sphere. Array dimension
 = 100.

AJR - array of spherical Bessel functions $j_n(x)$,
 $n = 1, \dots, N$ (complex x) values. Array dimension
= 100.

AM - complex index of refraction (m).

AMK - complex index of refraction (m) times wave
number (k).

CEX - complex time-variation factor $[\exp(i\omega t)]$.

Z - complex index of refraction (m) times Mie-size
parameter (x).

DP - array of spherical Bessel functions $\eta_n(x)$,
 $n = 1, \dots, N$ (real x) values or first
derivative values of associated Legendre functions
 $p_n^1(\theta)$, $n = 1, \dots, N$ (real θ). Array dimension = 100.

P - array of spherical Bessel functions $j_n(x)$,
 $n = 1, \dots, N$ (real x) values or values of
 $p_n^1(\theta)$, $n = 1, \dots, N$ (real θ). Array dimension = 100.

ALAMDA - wavelength of incident electric wave.

D1 - twice the radial distance to an interior point
of the sphere.

EO - intensity (field strength) of incident electric field.

PHI - azimuthal angle.

PIE - the number 3.14159265358973D0.

RKR - the reciprocal of Mie-size parameter.

STOPR - a test quantity for the termination of the recurrence
relation for determining the spherical Bessel functions
 $\eta_n(mx)$. It is equal in magnitude to 1.0D25.

SINTH - the value of $\sin\theta$.

THETA - the colatitude angle (θ) of an interior point of
the sphere.

NC - maximum order of spherical Bessel functions $\eta_n(mx)$
(complex mx) less two.

A single labeled common area, GAUSS, is used by the driver routine,
MAIN, and function subprogram, GAUSS3, for values of

D - diameter of the sphere.

M1 - number of points to be used in an m-point Gauss-Legendre quadrature rule to evaluate innermost integral of triple integral for total power.

M2 - number of points to be used in m-point Gauss-Legendre quadrature rule to evaluate middle integral of triple integral for total power.

In subroutine BESS, if variable N, the maximum order of spherical Bessel function $\eta_n(mx)$ (complex mx), tests ≤ 2 , the error message

PROCESS CAN NOT PROCEED SINCE NM2 = 0 FOR Z = ...

is printed out and the computer run is terminated.

In function subprogram, GAUSS3, an invalid value of one of the parameters M1, M2, or M3 institutes an error-message printout:

ERROR IN INTEGRATION CONTROLS. M1 = . . .

M2 = . . . M3 = . . .

A numerical value of zero for the triple integral, equation 54, is returned to driver routine, MAIN.

Input to program MIE is by keypunched cards. There are two basic input cards with structure and sequential order as follows:

Card No. 1 (control parameters)

Columns: 1-10 FREQ. Frequency in MHz. (E10.3)

11-20 EPS. Relative dielectric constant. (E10.3)

21-30 SIGMA1. Conductivity in mho/meter. (E10.3)

31-40 E0. Intensity (field strength) of incident electric field in volt/meter. (E10.3)

41-50 D. Diameter of sphere in cm. (E10.3)

51-60 TIME. Time in sec. (E10.3)

61-65 NOC. Number of cases. (I5)

66-70 IOPT = 0: Average power density and total power not computed; = 1: otherwise.

71-73 M1. Number of points for Gauss-Legendre quadrature rule applied to innermost integral of triple integral for total power deposited in the sphere.

74-76 M2. Same definition as M1 except application is to middle integral. (I3)

77-79 M3. Same definition as M1 except application is to outermost integral. (I3)

Cards Nos. 2-(NOC+1) (coordinate data)

Columns: 1-10 R. Radial spherical coordinate of interior point of sphere in cm. Range: $R > 0$. (E10.3)

11-20 THETAD. Colatitude spherical coordinate of interior point of sphere in deg. Range: $0 \leq \text{THETAD} \leq 180$. (E10.3)

21-30 PHID. Azimuthal spherical coordinate of interior point of sphere in deg. Range: $0 \leq \text{PHID} \leq 360$. (E10.3)

The last card of a single data set must be a termination card with the symbols /* punched in columns 1 and 2. Also program MIE can handle multiple data sets. Each data set [Cards 1-(NOC+1)] is stacked one behind the other, with the last card in the complete data deck a termination card.

Under option control IOPT=0, program output printouts consist of the title

DEPOSITION OF POWER INSIDE OF A HOMOGENEOUS SPHERE IMMERSED IN AN ELECTROMAGNETIC FIELD

followed by such information as

FREQUENCY = . . . MHZ

WAVELENGTH = . . . CM

CONDUCTIVITY = . . . MHO/M

RELATIVE DIELECTRIC CONST = . . .

DIAMETER = . . . CM

TIME = . . . SEC

REFRACTIVE INDEX = . . .

THE HIGHEST ORDER NEUMANN FUNCTION COMPUTED IS OF ORDER = . . .

SIZE PARAMETER = . . .

RATIO OF THE FIRST CALCULATION OF THE ZERO ORDER SPHERICAL
BESSEL FUNCTION OF ARGUMENT X TO THE VALUE OF $\sin(X)/X =$
. . .

RATIO OF THE FIRST CALCULATION OF THE ZERO ORDER SPHERICAL
BESSEL FUNCTION OF ARGUMENT Z TO THE VALUE OF $\sin(Z)/Z =$
. . .

INTERNAL POINT: RADIUS = . . . CM THETA = . . . DEG PHI = . . . DEG

ABSORBED POWER DENSITY = . . . WATTS/M**3

INTERNAL POINT: RADIUS = . . . CM THETA = . . . DEG PHI = . . . DEG

ABSORBED POWER DENSITY = . . . WATTS/M**3

For IOPT = 1, the printouts consist of the information set forth
above plus two additional statements:

TOTAL ABSORBED POWER = . . . WATTS

AVERAGE ABSORBED POWER DENSITY = . . . WATTS/M**3

Option IOPT = 0 with NOC = 0 produces a printout similar in
format to that presented on page 16 but without the data on the internal
points; whereas, option IOPT = 1 with NOC = 0 yields a like printout
plus total absorbed power and average absorbed power density results
formatted as in the above paragraph.

When using the IBM 360/65 system, the combined compilation and link
editing times of program MIE is about 0.30 minute. Execution time for
specific problems will be found in Appendix B.

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APPENDIX A

MATHEMATICAL ANALYSIS OF THE MIE SOLUTION

Consider four vector-valued functions (\underline{E} , \underline{H} , \underline{D} , and \underline{B}) defined on the Cartesian product ($R^3 \times R$) of space and time. Functions \underline{E} and \underline{H} denote the electric- and magnetic-field intensities; functions \underline{D} and \underline{B} denote electric displacement and magnetic induction respectively. Vector \underline{D} is also called electric induction by O'Rahilly (3, p. 35). These functions are related by Maxwell's equations

$$\nabla \times \underline{E} + \frac{1}{c} \frac{\partial \underline{B}}{\partial t} = 0, \quad (\text{A-1})$$

$$\nabla \times \underline{H} - \frac{1}{c} \frac{\partial \underline{D}}{\partial t} = \frac{4\pi\sigma \underline{E}}{c}, \quad (\text{A-2})$$

where σ is the conductivity of the medium at (x, y, z) at time t , and c is the velocity of light.

We assume the constitutive relations

$$\underline{B} = \mu \underline{H}, \quad (\text{A-3})$$

$$\underline{D} = \epsilon \underline{E}, \quad (\text{A-4})$$

and divide R^3 into two regions in each of which μ (magnetic permeability), ϵ (dielectric constant), and σ (conductivity) are constant functions.

We take the curl of both sides of equation A-1 and substitute into equation A-2 and have

$$\nabla(\nabla \cdot \underline{E}) - \nabla^2 \underline{E} + \frac{1}{c} \frac{\partial}{\partial t} (\nabla \times \underline{B}) = 0, \quad (\text{A-5})$$

which in turn implies that

$$\nabla(\nabla \cdot \underline{E}) - \nabla^2 \underline{E} + \frac{\mu}{c^2} \frac{\partial^2 \underline{D}}{\partial t^2} + \frac{4\pi\mu\sigma}{c^2} \frac{\partial \underline{E}}{\partial t} = 0. \quad (\text{A-6})$$

From equation A-6 we see that $\nabla \cdot \underline{E} = \rho/\epsilon$ implies

$$\nabla(\rho/\epsilon) - \nabla^2 \underline{E} + \frac{\mu\epsilon}{c^2} \frac{\partial^2 \underline{E}}{\partial t^2} + \frac{4\pi\mu\sigma}{c^2} \frac{\partial \underline{E}}{\partial t} = 0 \quad (\text{A-7})$$

or

$$\nabla(\nabla \cdot \underline{E}) - \nabla^2 \underline{E} + \frac{\mu\epsilon}{c^2} \frac{\partial^2 \underline{E}}{\partial t^2} + \frac{4\pi\mu\sigma}{c^2} \frac{\partial \underline{E}}{\partial t} = 0, \quad (\text{A-8})$$

where ρ is the charge density.

Let us assume that the \underline{E} field is

$$\underline{E} = \underline{E}_0 \exp(i\omega t). \quad (\text{A-9})$$

Then equation A-2 implies that

$$\nabla \times \underline{H} = \left(\frac{i\omega\epsilon}{c} \underline{E}_0 + \frac{4\pi\sigma}{c} \underline{E}_0 \right) \exp(i\omega t). \quad (\text{A-10})$$

From equation A-10 we see that if ϵ and σ are constant, then

$$\nabla \cdot (\nabla \times \underline{H}) = \left(\frac{i\omega\epsilon + 4\pi\sigma}{c} \right) \exp(i\omega t) \nabla \cdot \underline{E}_0 = 0. \quad (\text{A-11})$$

Thus

$$\nabla \cdot \underline{E} = 0 \quad (\text{A-12})$$

in a region where ϵ and σ are constant. Therefore, for time-harmonic waves, equation A-8 may be replaced by

$$-\nabla^2 \underline{E} + \frac{\mu\epsilon}{c^2} \frac{\partial^2 \underline{E}}{\partial t^2} + \frac{4\pi\mu\sigma}{c^2} \frac{\partial \underline{E}}{\partial t} = 0, \quad (\text{A-13})$$

in a region where ϵ and σ are constant. Substituting equation A-9 into equation A-13 we find that

$$\nabla^2 \underline{E} + \left(\frac{\mu \epsilon \omega^2}{c^2} - i \frac{4\pi \mu \sigma \omega}{c^2} \right) \underline{E} = 0. \quad (A-14)$$

Let

$$m^2 k^2 = \omega^2 \left[\frac{\mu \epsilon}{c^2} - i \frac{4\pi \mu \sigma}{\omega c^2} \right] = \left(\frac{\omega}{c} \right)^2 (\mu \epsilon - i \frac{4\pi \mu \sigma}{\omega}). \quad (A-15)$$

Then substituting equation A-15 into equation A-14, we have, with $k = \omega/c$ and $m^2 = \mu \epsilon - i 4\pi \mu \sigma / \omega$, the vector Helmholtz equation

$$\nabla^2 \underline{E} + m^2 k^2 \underline{E} = 0, \quad (A-16)$$

where k and m are the wave number and complex index of refraction, respectively, and remain as such in the rest of the paper.

Now the scalar-wave equation

$$\nabla^2 \psi + m^2 k^2 \psi = 0 \quad (A-17)$$

in spherical coordinates, a separable coordinate system for the equation, has the form

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \psi}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \psi}{\partial \phi^2} + m^2 k^2 \psi = 0. \quad (A-18)$$

Let

$$\psi = R(r) \Theta(\theta) \Phi(\phi). \quad (A-19)$$

Substituting equation A-19 into equation A-18 yields

$$\begin{aligned} \Theta \Phi \left[\frac{1}{r^2} \frac{\partial}{\partial r} (r^2 R') \right] + \frac{\Phi R}{r^2 \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta \Theta') \\ + \frac{\Theta R}{r^2 \sin^2 \theta} \Phi'' + m^2 k^2 R \Theta \Phi = 0. \end{aligned} \quad (A-20)$$

Expansion of terms in equation A-20 results in

$$\begin{aligned} \theta \phi [R'' + (2/r)R'] + \frac{\phi R}{r^2} \left(\frac{\cos \theta}{\sin \theta} \theta' + \theta'' \right) \\ + \frac{\theta R}{r^2 \sin^2 \theta} \phi'' + m^2 k^2 R \theta \phi = 0, \end{aligned} \quad (A-21)$$

and dividing all terms in equation A-21 by $R \theta \phi$, we infer that

$$\begin{aligned} \frac{1}{R} [R'' + (2/r)R'] + \left[\frac{1}{\theta r^2} \right] (\theta'' + \frac{\cos \theta}{\sin \theta} \theta') + \frac{1}{r^2 \sin^2 \theta} \phi'' \\ + m^2 k^2 = 0. \end{aligned} \quad (A-22)$$

Finally, multiplying all terms of equation A-22 by r^2 , we obtain the separated form of the equation

$$\begin{aligned} \left\{ \frac{r^2}{R} [R'' + (2/r)R'] + m^2 k^2 r^2 \right\} + \frac{1}{\theta} \left(\theta'' + \frac{\cos \theta}{\sin \theta} \theta' \right) \\ + \frac{1}{\sin^2 \theta} \phi'' = 0. \end{aligned} \quad (A-23)$$

We use the following lemma:

Lemma A-1. Let y be a Bessel function of order $(n + \frac{1}{2})$, then

$$x^2 y'' + x y' = [n(n+1) + \frac{1}{4} - x^2] y \quad (A-24)$$

along with

$$v^2 = n(n+1) \quad (A-25)$$

and

$$u = y / \sqrt{x} \quad (A-26)$$

implies that

$$x^2 u'' + 2xu' = (v^2 - x^2)u. \quad (A-27)$$

Proof of Lemma A-1. Equation A-24 is satisfied because

$(n+\frac{1}{2})^2 = n^2 + n + \frac{1}{4} = n(n+1) + \frac{1}{4}$ and by the definition of Bessel's equation (Whittaker [6, p. 38]). Differentiating u we find that

$$u' = y'/\sqrt{x} - (1/2)y/x^{3/2}, \quad (A-28)$$

$$\begin{aligned} u'' = y''/\sqrt{x} - (1/2)y'/x^{3/2} - (1/2)y'/x^{3/2} \\ + (3/4)y/x^{5/2}, \end{aligned} \quad (A-29)$$

or

$$u'' = y''/\sqrt{x} - y'/x^{3/2} + (3/4)y/x^{5/2}. \quad (A-30)$$

Now

$$\begin{aligned} x^2 u'' + 2xu' = x^{3/2}y'' - \sqrt{xy}' + (3/4)y/\sqrt{x} + 2\sqrt{xy}' \\ - y\sqrt{x} \end{aligned} \quad (A-31)$$

or

$$x^2 u'' + 2xu' = (1/\sqrt{x})[x^2 y'' + xy' - (1/4)y]. \quad (A-32)$$

From equations A-24 and A-32 it follows that

$$x^2 u'' + 2xu' = (1/\sqrt{x})\{[n(n+1) + (1/4)]y - (1/4)y-x^2 y\},$$

or

$$\begin{aligned}
x^2 u'' + 2xu' &= [n(n+1) - x^2](y/\sqrt{x}), \\
&= [n(n+1) - x^2]u.
\end{aligned}
\tag{A-33}$$

This completes the proof of Lemma A-1.

Now equation A-23 implies that there is a constant v^2 such that

$$(1/R)(r^2 R'' + 2rR') + m^2 k^2 r^2 = v^2 \tag{A-34}$$

or

$$r^2 R'' + 2rR' + (m^2 k^2 r^2 - v^2)R = 0. \tag{A-35}$$

Let us write

$$R(r) = z(mkr). \tag{A-36}$$

On substituting equation A-36 into equation A-35, we obtain

$$m^2 k^2 r^2 z''(mkr) + 2mkrz'(mkr) + (m^2 k^2 r^2 - v^2)z(mkr) = 0, \tag{A-37}$$

which according to Lemma A-1 is true provided that

$$z(mkr) = C \frac{Z_{n+\frac{1}{2}}(mkr)}{\sqrt{mkr}}, \tag{A-38}$$

where C is an arbitrary constant

$$v^2 = n(n+1), \tag{A-39}$$

and $Z_{n+\frac{1}{2}}$ is a Bessel function of order $n+\frac{1}{2}$.

Substituting equations A-36 and A-38 into equation A-23 results in the relation

$$n(n+1) + \frac{1}{\theta}(\theta'' + \frac{\cos \theta}{\sin \theta} \theta') + \frac{1}{\sin^2 \theta} \phi'' = 0. \tag{A-40}$$

Now multiplying all terms of equation A-40 by $\sin^2 \theta$, we find that

$$[n(n+1)\sin^2\theta + (1/\theta)(\sin^2\theta\theta'' + \sin\theta\cos\theta\theta')] + (1/\phi)\phi'' = 0. \quad (\text{A-41})$$

Let

$$(1/\phi)\phi'' = -\ell^2. \quad (\text{A-42})$$

Then

$$\phi'' + \ell^2\phi = 0 \quad (\text{A-43})$$

implies

$$\phi(\phi) = c_1\sin(\ell\phi) + c_2\cos(\ell\phi) \quad (\text{A-44})$$

for some constants, c_1 and c_2 . Substituting equation A-42 into equation A-41 yields

$$n(n+1)\sin^2\theta\theta + \sin^2\theta\theta'' + \sin\theta\cos\theta\theta' - \ell^2\theta = 0 \quad (\text{A-45})$$

or

$$(1-\cos^2\theta)\theta'' + \sin\theta\cos\theta\theta' + [n(n+1)(1-\cos^2\theta)-\ell^2]\theta = 0. \quad (\text{A-46})$$

Let us set

$$\theta = w(\cos\theta). \quad (\text{A-47})$$

Then

$$\theta' = -\sin\theta w'(\cos\theta), \quad (\text{A-48})$$

$$\theta'' = -\cos\theta w'(\cos\theta) + \sin^2\theta w''(\cos\theta). \quad (\text{A-49})$$

Substituting equations A-48 and A-49 into equation A-46, we obtain

$$(1-\cos^2\theta)[- \cos\theta w'(\cos\theta) + \sin^2\theta w''(\cos\theta)]$$

$$+ \cos\theta \sin^2\theta w'(\cos\theta) + [n(n+1)(1-\cos^2\theta) - \ell^2]w(\cos\theta) = 0$$

(A-50)

and dividing all terms of equation A-50 by $(1-\cos^2\theta)$ yields

$$(1-\cos^2\theta)w''(\cos\theta) - 2\cos\theta w'(\cos\theta)$$

$$+ [n(n+1) - \frac{\ell^2}{1-\cos^2\theta}] w(\cos\theta) = 0. \quad (A-51)$$

Now set $x = \cos\theta$ in equation A-51 and discover that $w(x)$ satisfies

$$(1-x^2)w''(x) - 2xw'(x) + [n(n+1) - \frac{\ell^2}{1-x^2}]w(x) = 0. \quad (A-52)$$

It can be shown that

$$w(x) = P_n^\ell(x) \quad (A-53)$$

is a solution of equation A-52, where $P_n^\ell(x)$ is a solution of the associated Legendre ordinary differential equation (Whittaker [6, p. 324]).

Hence

$$R\Theta\Phi = \frac{Z_{n+\frac{1}{2}}(mkr)}{\sqrt{mkr}} P_n^\ell(\cos\theta) [c_1 \cos(\ell\phi) + c_2 \sin(\ell\phi)], \quad (A-54)$$

where $Z_{n+\frac{1}{2}}$ denotes a Bessel function of order $n+\frac{1}{2}$. Observe that

$$\psi = R(r)\Theta(\theta)\Phi(\phi) \quad (A-55)$$

satisfies equation A-17.

The following vector functions:

$$\underline{L} = \underline{\nabla}\psi, \quad (\text{A-56})$$

$$\underline{M} = \underline{\nabla}\mathbf{x}(\underline{a}\psi) = \underline{\nabla}\psi\mathbf{x}\underline{a}, \quad (\text{A-57})$$

$$\underline{N} = (1/mk)\underline{\nabla}\mathbf{x}\underline{M}, \quad (\text{A-58})$$

in which \underline{a} is a constant vector and ψ is a spatial function, satisfy the same vector-wave equation since the components of these vectors are linear combinations of derivatives of solutions of the basic scalar Helmholtz wave equation--a linear partial differential equation with constant coefficients.

The validity of equation A-57 is based on the following proposition:

Proposition A-1. For every C^1 function ψ and every constant vector \underline{a} , we have

$$\underline{\nabla}\mathbf{x}(\underline{a}\psi) = \underline{\nabla}\psi\mathbf{x}\underline{a}. \quad (\text{A-59})$$

Proof of proposition A-1.

$$\begin{aligned} \underline{\nabla}\mathbf{x}\underline{a}\psi &= \begin{vmatrix} \underline{i} & \underline{j} & \underline{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ a_1\psi & a_2\psi & a_3\psi \end{vmatrix} \\ &= \underline{i}(a_3\frac{\partial\psi}{\partial y} - a_2\frac{\partial\psi}{\partial z}) - \underline{j}(a_3\frac{\partial\psi}{\partial x} - a_1\frac{\partial\psi}{\partial z}) \\ &\quad + \underline{k}(a_2\frac{\partial\psi}{\partial x} - a_1\frac{\partial\psi}{\partial y}), \end{aligned} \quad (\text{A-60})$$

$$\begin{aligned}
\underline{\nabla \psi \underline{x} \underline{a}} &= \begin{vmatrix} \underline{i} & \underline{j} & \underline{k} \\ \frac{\partial \psi}{\partial x} & \frac{\partial \psi}{\partial y} & \frac{\partial \psi}{\partial z} \\ a_1 & a_2 & a_3 \end{vmatrix} \\
&= \underline{i}(a_3 \frac{\partial \psi}{\partial y} - a_2 \frac{\partial \psi}{\partial z}) - \underline{j}(a_3 \frac{\partial \psi}{\partial x} - a_1 \frac{\partial \psi}{\partial z}) \\
&\quad + \underline{k}(a_2 \frac{\partial \psi}{\partial x} - a_1 \frac{\partial \psi}{\partial y}). \tag{A-61}
\end{aligned}$$

Since the right sides of equations A-60 and A-61 are identical, the proposition is proven.

The approach of Stratton (4, pp. 392-423) is to express the vector potential \underline{A} as a linear combination of vector functions \underline{L}_n , \underline{M}_n and \underline{N}_n where these depend on ψ_n , the nth spherical harmonic associated with the scalar Helmholtz wave equation.

Observe that

$$\underline{\nabla x}(\underline{r}\psi) = \underline{\nabla \psi} \underline{x} \underline{r} + \psi \underline{\nabla x} \underline{r}, \tag{A-62}$$

where \underline{r} is the position vector. To simplify equation A-62, we need to represent the basic operators in an orthogonal coordinate system.

Let (x,y,z) denote a point in Cartesian coordinates. Suppose that

$$x = x(u^1, u^2, u^3), \quad y = y(u^1, u^2, u^3), \quad z = z(u^1, u^2, u^3). \tag{A-63}$$

If the position vector, \underline{r} , is expressed as

$$\underline{r} = x\underline{i} + y\underline{j} + z\underline{k} \quad (\underline{i}, \underline{j}, \underline{k} \text{ unit base vectors}), \tag{A-64}$$

then in Cartesian coordinates

$$\underline{\nabla} \psi \underline{xr} = \begin{vmatrix} \underline{i} & \underline{j} & \underline{k} \\ \psi_x & \psi_y & \psi_z \\ x & y & z \end{vmatrix}$$

$$= \underline{i}(z\psi_y - y\psi_z) - \underline{j}(z\psi_x - x\psi_z) + \underline{k}(y\psi_x - x\psi_y). \quad (\text{A-65})$$

It is therefore natural to define operators L_x , L_y , and L_z by the rules

$$L_x \psi = z \frac{\partial \psi}{\partial y} - y \frac{\partial \psi}{\partial z}, \quad (\text{A-66})$$

$$L_y \psi = x \frac{\partial \psi}{\partial z} - z \frac{\partial \psi}{\partial x}, \quad (\text{A-67})$$

$$L_z \psi = y \frac{\partial \psi}{\partial x} - x \frac{\partial \psi}{\partial y}. \quad (\text{A-68})$$

Note that

$$\begin{aligned} L_x \nabla^2 \psi &= z \left(\frac{\partial^3 \psi}{\partial x^2 \partial y} + \frac{\partial^3 \psi}{\partial y^3} + \frac{\partial^3 \psi}{\partial z^2 \partial y} \right) - y \left(\frac{\partial^3 \psi}{\partial z \partial x^2} + \frac{\partial^3 \psi}{\partial z \partial y^2} + \frac{\partial^3 \psi}{\partial z^3} \right) \\ &= \left[\left(\frac{\partial}{\partial x} \right)^2 + \left(\frac{\partial}{\partial y} \right)^2 + \left(\frac{\partial}{\partial z} \right)^2 \right] \left(z \frac{\partial \psi}{\partial y} - y \frac{\partial \psi}{\partial z} \right). \end{aligned} \quad (\text{A-69})$$

That is to say, $L_x \nabla^2 = \nabla^2 L_x$. Moreover, we have the following theorem:

Theorem A-1. For all C^3 functions ψ , we have

$$\nabla^2 L_x \psi = L_x \nabla^2 \psi, \quad (\text{A-70})$$

$$\nabla^2 L_y \psi = L_y \nabla^2 \psi, \quad (\text{A-71})$$

$$\nabla^2 L_z \psi = L_z \nabla^2 \psi. \quad (\text{A-72})$$

Hence, if

$$\nabla^2 \psi = -m^2 k^2 \psi, \quad (\text{A-73})$$

the vector function

$$\begin{aligned} \underline{V} &= \underline{iL}_x \psi + \underline{jL}_y \psi + \underline{kL}_z \psi \\ &= \underline{\nabla \psi \times r} \end{aligned} \quad (\text{A-74})$$

satisfies

$$\underline{\nabla}^2 \underline{V} = -m^2 k^2 \underline{V}. \quad (\text{A-75})$$

Proof of Theorem A-1. If \underline{V} is defined by equation A-74, then equations A-70 through A-73 imply that

$$\begin{aligned} \underline{\nabla}^2 \underline{V} &= \underline{i} \underline{\nabla}^2 L_x \psi + \underline{j} \underline{\nabla}^2 L_y \psi + \underline{k} \underline{\nabla}^2 L_z \psi \\ &= \underline{iL}_x \underline{\nabla}^2 \psi + \underline{jL}_y \underline{\nabla}^2 \psi + \underline{kL}_z \underline{\nabla}^2 \psi \\ &= \underline{iL}_x (-m^2 k^2 \psi) + \underline{jL}_y (-m^2 k^2 \psi) + \underline{kL}_z (-m^2 k^2 \psi) \\ &= -m^2 k^2 \underline{V}. \end{aligned} \quad (\text{A-76})$$

This establishes Theorem A-1.

If u and v are two solutions of the scalar-wave equation, then Maxwell's equations are satisfied by

$$\underline{E} = \underline{M}_v + i \underline{N}_u \quad (i \text{ is the complex unit}). \quad (\text{A-77})$$

Equation A-1 implies

$$\underline{\nabla \times E} + \frac{i\omega}{c} \underline{B} = 0. \quad (\text{A-78})$$

Thus

$$\underline{H} = - \frac{c}{i\mu\omega} \nabla \times \underline{E} = i \left(\frac{c}{\mu\omega} \right) \nabla \times \underline{E} . \quad (A-79)$$

Furthermore, in view of equations A-58 and A-77

$$\underline{H} = \left(\frac{ic}{\mu\omega} \right) m \underline{N}_V + \left(\frac{ic}{\mu\omega} \right) [\nabla \times (\nabla \times \underline{M}_U) / mk] i \quad (A-80)$$

or, equivalently,

$$\underline{H} = \frac{imkc}{\mu\omega} \underline{N}_V + \left(\frac{ic}{\mu\omega mk} \right) [\nabla (\nabla \cdot \underline{M}_U) - \nabla^2 \underline{M}_U] i . \quad (A-81)$$

The definition of \underline{M}_U by the rule

$$\underline{M}_U = \nabla \times \underline{r} u, \quad (A-82)$$

the vector identity (the divergence of a curl is zero), and equations A-75 or A-76, imply that

$$\underline{H} = \frac{imkc}{\mu\omega} \underline{N}_V + \frac{imkc}{\mu\omega} \underline{M}_U i \quad (A-83)$$

or

$$\underline{H} = \frac{mkc}{\mu\omega} (-\underline{M}_U + i \underline{N}_V) . \quad (A-84)$$

In the Gaussian units for a nonmagnetic body, $\mu = 1$. Thus, since $k = \omega/c$, we have

$$\underline{H} = m(-\underline{M}_U + i \underline{N}_V) . \quad (A-85)$$

In the spherical coordinate system with unit base vectors, \underline{e}_r , \underline{e}_θ , and \underline{e}_ϕ , vector \underline{A} can be represented as

$$\underline{A} = A_r \underline{e}_r + A_\theta \underline{e}_\theta + A_\phi \underline{e}_\phi . \quad (\text{A-85.1})$$

Then

$$\begin{aligned} \underline{\nabla \times A} &= \frac{1}{r^2 \sin \theta} \{ [\frac{\partial}{\partial \theta}(r \sin \theta A_\phi) - \frac{\partial}{\partial \phi}(r A_\theta)] \underline{e}_r \\ &+ [\frac{\partial A_r}{\partial \phi} - \frac{\partial}{\partial r}(r \sin \theta A_\phi)] r \underline{e}_\theta \\ &+ [\frac{\partial}{\partial r}(r A_\theta) - \frac{\partial A_r}{\partial \theta}] r \sin \theta \underline{e}_\phi \} . \end{aligned} \quad (\text{A-86})$$

We take $A_r = r\psi$, $A_\theta = 0$, and $A_\phi = 0$, and observe that

$$\underline{M}_\psi = M_r \underline{e}_r + M_\theta \underline{e}_\theta + M_\phi \underline{e}_\phi \quad (\text{A-87})$$

implies that

$$M_r = 0, \quad (\text{A-88})$$

$$M_\theta = \frac{1}{r \sin \theta} \frac{\partial}{\partial \phi}(r\psi), \quad (\text{A-89})$$

$$M_\phi = -\frac{1}{r} \frac{\partial}{\partial \theta}(r\psi). \quad (\text{A-90})$$

Furthermore, the functions N_ψ satisfy the relation

$$\underline{N}_\psi = \frac{1}{mk} \underline{\nabla \times M}_\psi , \quad (\text{A-91})$$

which implies that if

$$N_{\psi} = N_{\frac{e}{r-r}} + N_{\frac{e}{\theta-\theta}} + N_{\frac{e}{\phi-\phi}}, \quad (\text{A-92})$$

then

$$N_r = \frac{1}{mk} \left\{ - \left(\frac{1}{r^2 \sin \theta} \right) \frac{\partial}{\partial \theta} [\sin \theta \frac{\partial}{\partial \theta} (r\psi)] - \left(\frac{1}{r^2 \sin^2 \theta} \right) \frac{\partial^2}{\partial \phi^2} (r\psi) \right\}, \quad (\text{A-93})$$

$$N_{\theta} = \frac{1}{mk} \left\{ - \frac{\partial}{\partial r} [r \sin \theta \left(- \frac{1}{r} \frac{\partial}{\partial \theta} (r\psi) \right)] r \right\} \left(\frac{1}{r^2 \sin \theta} \right)$$

$$= \frac{1}{mk} \left[\frac{1}{r} \frac{\partial^2}{\partial r \partial \theta} (r\psi) \right], \quad (\text{A-94})$$

$$N_{\phi} = \frac{1}{mkr^2 \sin \theta} \left\{ \frac{\partial}{\partial r} \left[\frac{r}{r \sin \theta} \frac{\partial}{\partial \phi} (r\psi) \right] \right\} r \sin \theta, \quad (\text{A-95})$$

Hence

$$N_{\phi} = \frac{1}{mkr \sin \theta} \frac{\partial^2}{\partial r \partial \phi} (r\psi). \quad (\text{A-96})$$

Now equation A-93 may be simplified by using the fact that ψ satisfies equation A-18. We have

$$N_r = \frac{1}{mk} \left[\frac{1}{r} \frac{\partial}{\partial r} (r^2 \frac{\partial \psi}{\partial r}) + m^2 k^2 r \psi \right]. \quad (\text{A-97})$$

Since

$$\frac{\partial}{\partial r} (r\psi) = r \frac{\partial \psi}{\partial r} + \psi, \quad (\text{A-98})$$

partial differentiation with respect to r yields

$$\frac{\partial^2}{\partial r^2}(r\psi) = r \frac{\partial^2 \psi}{\partial r^2} + 2 \frac{\partial \psi}{\partial r} = \frac{1}{r} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \psi}{\partial r} \right) . \quad (\text{A-99})$$

Substituting equation A-99 into equation A-97, we see that

$$N_r = \frac{1}{mk} \left[-\frac{\partial^2}{\partial r^2}(r\psi) + m^2 k^2 r\psi \right] . \quad (\text{A-100})$$

The incident wave is described by the vectors

$$\underline{E} = \underline{a}_x E_o \exp[-i(kz - \omega t)] , \quad (\text{A-101})$$

$$\underline{H} = \underline{a}_y E_o \exp[-i(kz - \omega t)] , \quad (\text{A-102})$$

where vectors \underline{a}_x and \underline{a}_y are unit vectors directed along the x, y -axes. Vectors \underline{E} and \underline{H} can be expressed also in the form of equations A-77 and A-85 with scalar functions

$$u = E_o \exp(i\omega t) \cos\phi \sum_{n=1}^{\infty} (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos\theta) j_n(kr) , \quad (\text{A-103})$$

$$v = E_o \exp(i\omega t) \sin\phi \sum_{n=1}^{\infty} (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos\theta) j_n(kr) . \quad (\text{A-104})$$

The representation of vector \underline{E} is given by

$$\underline{E} = \underline{M}_v + i \underline{N}_u . \quad (\text{A-105})$$

The terms

$$(\underline{M}_v)_\theta = \frac{1}{\sin\theta} \frac{\partial v}{\partial \phi} = F_\theta(\theta, \phi) v , \quad (\text{A-106})$$

$$(\underline{N}_u)_\theta = \frac{1}{mk} \frac{1}{r} \frac{\partial^2}{\partial r \partial \theta} (ru) = G_\theta(\theta, \phi) \frac{\partial}{\partial r} (ru) \quad (A-107)$$

make up E_θ . Similarly, for E_ϕ

$$(\underline{M}_v)_\phi = - \frac{\partial v}{\partial \theta} = F_\phi(\theta, \phi) v, \quad (A-108)$$

$$(\underline{N}_u)_\phi = \frac{1}{mkr \sin \theta} \frac{\partial^2}{\partial r \partial \phi} (ru) = G_\phi(\theta, \phi) \frac{\partial}{\partial r} (ru) \quad (A-109)$$

This does not remain true when in equations A-103 and A-104 the $(-i)^n$ is replaced by $(-i)^n a_n$ and $(-i)^n b_n$, respectively, where the $(a_1, a_2, \dots, a_n, \dots)$ and $(b_1, b_2, \dots, b_n, \dots)$ are members of the sequence space appropriate to the functions being represented. We see by differentiation in the tangential directions that if

$\frac{1}{m} \frac{\partial}{\partial r} (ru)$ and v are continuous in r for each θ and ϕ , then

E_θ and E_ϕ are continuous across the boundary of the sphere.

Now we repeat this argument for the \underline{H} vector. The \underline{H} vector is defined by

$$\underline{H} = m(-\underline{M}_u + i\underline{N}_v). \quad (A-110)$$

To get the tangential components of \underline{H} , simply interchange u and v in the preceding argument. Thus, the continuity of tangential \underline{H} is assured if $\frac{\partial(rv)}{\partial r}$ and mu are continuous in r for each θ and ϕ .

We introduce a new set of functions which differ from spherical Bessel functions by an additional factor, z . These functions are defined by

$$\psi_n(z) = z j_n(z) = (\pi z/2)^{1/2} J_{n+1/2}(z), \quad (A-111)$$

$$\chi_n(z) = -z \eta_n(z) = -(\pi z/2)^{1/2} N_{n+1/2}(z), \quad (A-112)$$

$$\zeta_n(z) = zh_n^{(2)}(z) = (\pi z/2)^{1/2} h_{n+1/2}^{(2)}(z), \quad (\text{A-113})$$

where

$$J_\nu(z) = \left[\sum_{k=0}^{\infty} \frac{(-1)^k (z/2)^{2k}}{\Gamma(k+1)\Gamma(k+\nu+1)} \right] (z/2)^\nu \quad (\text{A-114})$$

is holomorphic for integral ν and converges for all ν if the appropriate definition of $(z/2)^\nu$ is given and

$$Y_\nu(z) = \frac{J_\nu(z)\cos(\nu\pi) - J_{-\nu}(z)}{\sin(\nu\pi)}, \quad (\text{A-115})$$

$$\begin{aligned} H_\nu^{(2)}(z) &= J_\nu(z) - i Y_\nu(z) \\ &= \frac{J_{-\nu}(z) - \exp(i\nu\pi)J_\nu(z)}{-i \sin(\nu\pi)}, \end{aligned} \quad (\text{A-116})$$

$$N_\nu(z) = Y_\nu(z). \quad (\text{A-117})$$

(The functions N_ν are called Bessel functions of the second kind, Neumann functions, or Weber functions.)

An outgoing scattered wave can be obtained by applying a linear operator to the ordered pair (u, ν) , where

$$u = E_0 \exp(i\omega t) \cos\phi \sum_{n=1}^{\infty} -a_n (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos\theta) h_n^{(2)}(kr), \quad (\text{A-118})$$

$$\nu = E_0 \exp(i\omega t) \sin\phi \sum_{n=1}^{\infty} -b_n (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos\theta) h_n^{(2)}(kr). \quad (\text{A-119})$$

Similarly, an induced wave can be obtained by applying a linear operator to (u, v) , where

$$u = E_0 \exp(i\omega t) \cos\phi \sum_{n=1}^{\infty} m c_n (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos\theta) j_n(mkr), \quad (A-120)$$

$$v = E_0 \exp(i\omega t) \sin\phi \sum_{n=1}^{\infty} m d_n (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos\theta) j_n(mkr). \quad (A-121)$$

The boundary conditions are defined by

$$\underline{E}_{\text{tangent}}^{\text{incident}} + \underline{E}_{\text{tangent}}^{\text{scattered}} = \underline{E}_{\text{tangent}}^{\text{inside}} \quad (A-122)$$

and

$$\underline{H}_{\text{tangent}}^{\text{incident}} + \underline{H}_{\text{tangent}}^{\text{scattered}} = \underline{H}_{\text{tangent}}^{\text{inside}}. \quad (A-123)$$

Since the tangent space of the sphere is two-dimensional, we have four equations in four unknowns. These are assured by the continuity of

μ , $\frac{1}{m} \frac{\partial(ru)}{\partial r}$, v , and $\frac{\partial(rv)}{\partial r}$. Thus, equations A-103 and A-104,

equations A-118 through A-123, and the continuity relations imply that the relations that must be satisfied are

$$\begin{matrix} (\mu) \\ (A-103) \end{matrix} + \begin{matrix} (\mu) \\ (A-118) \end{matrix} = \begin{matrix} (\mu) \\ (A-120) \end{matrix}, \quad (A-124)$$

$$\begin{matrix} (\frac{1}{m} \frac{\partial(ru)}{\partial r}) \\ (A-103) \end{matrix} + \begin{matrix} (\frac{1}{m} \frac{\partial(ru)}{\partial r}) \\ (A-118) \end{matrix} = \begin{matrix} (\frac{1}{m} \frac{\partial(ru)}{\partial r}) \\ (A-120) \end{matrix}, \quad (A-125)$$

$$v_{(A-104)} + v_{(A-119)} = v_{(A-121)}, \quad (A-126)$$

$$\begin{matrix} (\frac{\partial(rv)}{\partial r}) \\ (A-104) \end{matrix} + \begin{matrix} (\frac{\partial(rv)}{\partial r}) \\ (A-119) \end{matrix} = \begin{matrix} (\frac{\partial(rv)}{\partial r}) \\ (A-121) \end{matrix}. \quad (A-127)$$

In equations A-124 and A-125, the m denotes the refractive index of the propagating medium. Outside the sphere, $m = m_o = 1$.

Equation A-124 with $r = a$ yields

$$\begin{aligned}
 & m_o E_o \exp(i\omega t) \cos\phi \sum_{n=1}^{\infty} (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos\theta) j_n(ka) \\
 & + m_o E_o \exp(i\omega t) \cos\phi \sum_{n=1}^{\infty} -a_n (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos\theta) h_n^{(2)}(ka) \\
 & = m E_o \exp(i\omega t) \cos\phi \sum_{n=1}^{\infty} m c_n (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos\theta) j_n(mka). \quad (A-128)
 \end{aligned}$$

Multiplying all terms of equation A-128 by

$$x = ka, \quad (A-129)$$

setting

$$y = mka = mx, \quad (A-130)$$

and using the fact that

$$m_o = 1, \quad (A-131)$$

we finally arrive at

$$\begin{aligned}
 x j_n(x) - a_n x h_n^{(2)}(x) &= m^2 c_n x j_n(mka) \\
 &= m c_n y j_n(y) \\
 &= m c_n \psi_n(y). \quad (A-132)
 \end{aligned}$$

Equations A-111 through A-113 and A-132 yield

$$\psi_n(x) - a_n \zeta_n(x) = m c_n \psi_n(y). \quad (A-133)$$

Next we make use of the continuity of $\frac{1}{m} \frac{\partial(ru)}{\partial r}$. Using equations A-103, A-118, A-120, and A-125 results in

$$\begin{aligned} & \frac{1}{m_o k} E_o \exp(i\omega t) \cos \phi \sum_{n=1}^{\infty} (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos \theta) \frac{\partial}{\partial r} (kr j_n(kr)) \\ & + \frac{1}{m_o k} E_o \exp(i\omega t) \cos \phi \sum_{n=1}^{\infty} (-a_n (-i)^n) \frac{2n+1}{n(n+1)} P_n^1(\cos \theta) \frac{\partial}{\partial r} (kr h_n^{(2)}(kr)) \\ & = \frac{1}{mk} E_o \exp(i\omega t) \cos \phi \sum_{n=1}^{\infty} mc_n (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos \theta) \frac{\partial}{\partial r} (mkr j_n(mkr)). \quad (A-134) \end{aligned}$$

Thus, using equations A-111 through A-113 and the orthogonality of the $P_n^1(\cos \theta)$, we arrive at

$$\psi'_n(x) - a_n \zeta'_n(x) = c_n \psi'_n(y). \quad (A-135)$$

Here we have used the fact that $y = mkr$ and that

$$\frac{\partial}{\partial r} [mkr j_n(mkr)] = mk \frac{\partial}{\partial y} [y j_n(y)]. \quad (A-136)$$

The continuity of v is expressed in equation A-126, and use of A-104, A-119, and A-121 shows that

$$\begin{aligned} & E_o \exp(i\omega t) \sin \phi \sum_{n=1}^{\infty} (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos \theta) j_n(kr) \\ & + E_o \exp(i\omega t) \sin \phi \sum_{n=1}^{\infty} -b_n (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos \theta) h_n^{(2)}(kr) \\ & = E_o \exp(i\omega t) \sin \phi \sum_{n=1}^{\infty} md_n (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos \theta) j_n(mkr). \quad (A-137) \end{aligned}$$

Thus, multiplying both sides of equation A-137 by kr and using equations A-111 through A-113, A-129, and A-130, we have

$$\psi_n(x) - b_n \zeta_n(x) = d_n \psi_n(y). \quad (\text{A-138})$$

The continuity of $\frac{\partial(rv)}{\partial r}$ is expressed by the equation A-127 and this, in view of equations A-104, A-119, and A-121, implies that

$$\begin{aligned} & E_o \exp(i\omega t) \sin\phi \sum_{n=1}^{\infty} (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos\theta) \frac{\partial}{\partial r} [r j_n(kr)] \\ & + E_o \exp(i\omega t) \sin\phi \sum_{n=1}^{\infty} -b_n (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos\theta) \frac{\partial}{\partial r} [r h_n^{(2)}(kr)] \\ & = E_o \exp(i\omega t) \sin\phi \sum_{n=1}^{\infty} m d_n (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos\theta) \frac{\partial}{\partial r} [r j_n(mkr)]. \end{aligned} \quad (\text{A-139})$$

Using equations A-129, A-130, A-136, and A-139, and the orthogonality of the $P_n^1(\cos\theta)$, we deduce that

$$\psi'_n(x) - b_n \zeta'_n(x) = m d_n \psi'_n(y). \quad (\text{A-140})$$

Dividing the right side of equation A-133 by the right side of equation A-135, we obtain

$$\frac{\psi_n(x) - a_n \zeta_n(x)}{\psi'_n(x) - a_n \zeta'_n(x)} = \frac{m \psi_n(y)}{\psi'_n(y)} \quad (\text{A-141})$$

which implies that

$$\frac{\psi_n(x) \psi'_n(y) - m \psi_n(y) \psi'_n(x)}{\psi'_n(y) \zeta_n(x) - m \psi_n(y) \zeta'_n(x)} = a_n. \quad (\text{A-142})$$

Similarly, equations A-138 and A-140 imply that

$$\frac{m\psi'_n(y)\psi_n(x) - \psi_n(y)\psi'_n(x)}{m\psi'_n(y)\zeta_n(x) - \psi_n(y)\zeta'_n(x)} = b_n. \quad (\text{A-143})$$

Observe that

$$\frac{d}{dx}[\psi'_n(x)\zeta_n(x) - \psi_n(x)\zeta'_n(x)] = \psi''_n(x)\zeta_n(x) - \psi_n(x)\zeta''_n(x). \quad (\text{A-144})$$

Now

$$\psi''_n(x) = xj''_n(x) + 2j'_n(x) \quad (\text{A-145})$$

implies that

$$\begin{aligned} x\psi''_n(x) &= x^2j''_n(x) + 2xj'_n(x) \\ &= [n(n+1) - x^2]j_n(x). \end{aligned} \quad (\text{A-146})$$

Similarly,

$$x\zeta''_n(x) = [n(n+1) - x^2]h^{(2)}_n(x). \quad (\text{A-147})$$

Thus, in view of definitions A-111 and A-113

$$x \frac{d}{dx} [\psi'_n(x)\zeta_n(x) - \psi_n(x)\zeta'_n(x)] = 0, \quad (\text{A-148})$$

which implies that

$$\psi'_n(x)\zeta_n(x) - \psi_n(x)\zeta'_n(x) = c, \quad (\text{A-149})$$

where c is independent of x .

Lemma A-2. For all complex z , we have

$$\psi'_n(z)\zeta_n(z) - \psi_n(z)\zeta'_n(z) = i. \quad (\text{A-150})$$

Proposition A-2. For every positive integer n and for x and y defined by equations A-129 and A-130, we have

$$c_n = \frac{i}{\psi'_n(y)\zeta_n(x) - m\psi_n(y)\zeta'_n(x)}, \quad (\text{A-151})$$

$$d_n = \frac{i}{m\psi'_n(y) - \psi_n(y)\zeta'_n(x)}. \quad (\text{A-152})$$

Proof of Proposition A-2. From equation A-133 it follows that

$$\psi_n(x) - mc_n\psi_n(y) = a_n\zeta_n(x), \quad (\text{A-153})$$

and from equation A-135 it follows that

$$\psi'_n(x) - c_n\psi'_n(y) = a_n\zeta'_n(x). \quad (\text{A-154})$$

Dividing the right side of equation A-153 by the right side of equation A-154, we deduce that

$$\frac{\psi_n(x) - mc_n\psi_n(y)}{\psi'_n(x) - c_n\psi'_n(y)} = \frac{\zeta_n(x)}{\zeta'_n(x)}, \quad (\text{A-155})$$

or, equivalently,

$$\begin{aligned}\psi_n(x)\zeta'_n(x) - mc_n\psi_n(y)\zeta'_n(x) \\ = \zeta_n(x)\psi'_n(x) - c_n\psi'_n(y)\zeta_n(x) .\end{aligned}\quad (A-156)$$

Thus

$$c_n = \frac{\psi'_n(x)\zeta_n(x) - \psi_n(x)\zeta'_n(x)}{\psi'_n(y)\zeta_n(x) - m\psi_n(y)\zeta'_n(x)} . \quad (A-157)$$

By Lemma A-2, equation A-157 implies equation A-151.

Equation A-138 leads to

$$\psi_n(x) - d_n\psi_n(y) = b_n\zeta_n(x) , \quad (A-158)$$

and equation A-140 leads to

$$\psi'_n(x) - md_n\psi'_n(y) = b_n\zeta'_n(x) . \quad (A-159)$$

Division of the right side of equation A-158 by the right side of equation A-159 results in the expression

$$\frac{\psi_n(x) - d_n\psi_n(y)}{\psi'_n(x) - md_n\psi'_n(y)} = \frac{\zeta_n(x)}{\zeta'_n(x)} . \quad (A-160)$$

Thus

$$d_n = \frac{\psi'_n(x)\zeta_n(x) - \psi_n(x)\zeta'_n(x)}{m\psi'_n(y)\zeta_n(x) - \psi_n(y)\zeta'_n(x)} . \quad (A-161)$$

We obtain equation A-152 by using equation A-161 and Lemma A-2. This completes the proof of Proposition A-2.

Proof of Lemma A-2. Observe that

$$J_\nu(z) = \left[\sum_{k=0}^{\infty} \frac{(-1)^k (z/2)^k}{\Gamma(k+1)\Gamma(k+\nu+1)} \right] (z/2)^\nu. \quad (\text{A-162})$$

Thus, definition A-111 implies that

$$\begin{aligned} \psi_n(z) &= (\pi z/2)^{\frac{1}{2}} \sum_{k=0}^{\infty} \frac{(-1)^k (z/2)^k}{\Gamma(k+1)\Gamma(k+n+\frac{1}{2}+1)} (z/2)^n (z/2)^{\frac{1}{2}} \\ &= \sqrt{\pi} \sum_{k=0}^{\infty} \frac{(-1)^k (z/2)^k}{\Gamma(k+1)\Gamma(k+n+3/2)} (z/2)^{n+1}, \end{aligned} \quad (\text{A-163})$$

and

$$\psi'_n(z) = \sqrt{\pi} \sum_{k=0}^{\infty} \frac{(-1)^k (n+k+1) (z/2)^k z^{n+1}}{\Gamma(k+1)\Gamma(k+n+3/2)}. \quad (\text{A-164})$$

In addition, definitions A-113 and A-116 imply that

$$\zeta_n(z) = (\pi z/2)^{\frac{1}{2}} \left\{ J_{n+\frac{1}{2}}(z) - i \left[\frac{J_{n+\frac{1}{2}}(z) \cos((n+\frac{1}{2})\pi) - J_{-(n+\frac{1}{2})}(z)}{\sin((n+\frac{1}{2})\pi)} \right] \right\} \quad (\text{A-165})$$

or

$$\zeta_n(z) = (\pi z/2)^{\frac{1}{2}} [J_{n+\frac{1}{2}}(z) - i(-1)^{n+1} J_{-(n+\frac{1}{2})}(z)]. \quad (\text{A-166})$$

Hence, equation A-166 implies that

$$\begin{aligned} \zeta_n(z) = & \sqrt{\pi} \sum_{k=0}^{\infty} \frac{(-1)^k (z/2)^k}{\Gamma(k+1) \Gamma(k+n+3/2)} (z/2)^{n+1} \\ & + i(-1)^n \sqrt{\pi} \sum_{k=0}^{\infty} \frac{(-1)^k (z/2)^k}{\Gamma(k+1) \Gamma[k-(n+\frac{1}{2})+1]} (z/2)^{-n}, \end{aligned} \quad (\text{A-167})$$

$$\begin{aligned} \zeta'_n(z) = & \sqrt{\pi} \sum_{k=0}^{\infty} \frac{(-1)^k (n+k+1) (z/2)^k z^{n/2} 2^{n+1}}{\Gamma(k+1) \Gamma(k+n+3/2)} \\ & + i(-1)^n \sqrt{\pi} \sum_{k=0}^{\infty} \frac{(-1)^k (k-n) 2^n (z/2)^k z^{-n-1}}{\Gamma(k+1) \Gamma[k-(n+\frac{1}{2})+1]}. \end{aligned} \quad (\text{A-168})$$

The nonzero terms are

$$\psi_n(z) \zeta'_n(z) \Big|_{z=0} = \frac{i(-1)^n \pi 2^n (-n) (\frac{1}{2})^{n+1}}{\Gamma(1) \Gamma(\frac{1}{2}-n) \Gamma(1) \Gamma(n+3/2)}, \quad (\text{A-169})$$

$$\psi'_n(z) \zeta_n(z) \Big|_{z=0} = \frac{i(-1)^n \pi (\frac{1}{2}) (n+1)}{\Gamma(1) \Gamma(n+3/2) \Gamma(\frac{1}{2}-n) \Gamma(1)}. \quad (\text{A-170})$$

Finally, we observe that

$$\begin{aligned} \zeta_n(z) \psi'_n(z) - \psi_n(z) \zeta'_n(z) &= A_n \\ &= \frac{i(-1)^n \pi (n+\frac{1}{2})}{\Gamma^2(1) \Gamma(n+3/2) \Gamma(\frac{1}{2}-n)}. \end{aligned} \quad (\text{A-171})$$

To complete the computation, we observe that

$$\Gamma(-\frac{1}{2}) = -2 \Gamma(\frac{1}{2}) \quad (\text{A-172})$$

since

$$\Gamma(-\frac{1}{2}+1) = (-\frac{1}{2}) \Gamma(-\frac{1}{2}) \quad (\text{A-173})$$

follows from the known relation

$$\Gamma(z+1) = z \Gamma(z) . \quad (\text{A-174})$$

Furthermore, $n = 1$ and equation A-171 implies that

$$A_1 = \frac{i(-1)(3/2)}{\Gamma^2(1)\Gamma(1+3/2)\Gamma(-\frac{1}{2})} . \quad (\text{A-175})$$

Application of equation A-174 to equation A-175 results in

$$\begin{aligned} A_1 &= \frac{-i(3/2)\pi}{(3/2)(\frac{1}{2})\Gamma(\frac{1}{2})\Gamma(\frac{1}{2})(-2)} \\ &= \frac{i\pi}{\Gamma^2(\frac{1}{2})} . \end{aligned} \quad (\text{A-176})$$

Now we see that for $\text{Re } z > 0$,

$$\Gamma(z) = \int_0^\infty t^{z-1} e^{-t} dt \quad (\text{A-177})$$

since

$$\begin{aligned}
 \Gamma(z+1) &= \int_0^{\infty} t^z e^{-t} dt \\
 &= - \int_0^{\infty} t^z de^{-t} \\
 &= -t^z e^{-t} \Big|_0^{\infty} + \int_0^{\infty} z t^{z-1} e^{-t} dt \\
 &= z \Gamma(z) .
 \end{aligned}
 \tag{A-178}$$

Observe that

$$\begin{aligned}
 \Gamma(1) &= \int_0^{\infty} e^{-t} dt \\
 &= 1,
 \end{aligned}
 \tag{A-179}$$

and

$$\Gamma(\tfrac{1}{2}) = \int_0^{\infty} t^{-\frac{1}{2}} e^{-t} dt.
 \tag{A-180}$$

Substituting $t = s^2$ in equation A-180, we see that

$$\begin{aligned}
 \Gamma(\tfrac{1}{2}) &= \int_0^{\infty} (1/s) e^{-s^2} 2s ds \\
 &= 2 \int_0^{\infty} e^{-s^2} ds \\
 &= \sqrt{\pi}
 \end{aligned}
 \tag{A-181}$$

since

$$I = \int_0^{\infty} e^{-s^2} ds \quad (A-182)$$

implies that

$$\begin{aligned} I^2 &= \int_0^{\pi/2} \int_0^{\infty} e^{-r^2} r dr d\theta \\ &= \pi/4 \end{aligned} \quad (A-183)$$

or

$$I = \sqrt{\pi}/2 . \quad (A-184)$$

Thus, equations A-176 and A-181 through A-183 imply that

$$A_1 = i . \quad (A-185)$$

We claim that $A_n = i$ for all n . Our claim is established by mathematical induction on n . Suppose that $n > 1$ and that $A_{n-1} = i$.

Then

$$\Gamma(\frac{1}{2}-n) = \frac{\Gamma(3/2-n)}{(\frac{1}{2}-n)} \quad (A-186)$$

and

$$\Gamma(n+3/2) = (n+\frac{1}{2})\Gamma(n+\frac{1}{2}) \quad (A-187)$$

imply that

$$A_n = \frac{i(-1)^n \pi(n+\frac{1}{2})}{(n+\frac{1}{2})\Gamma(n+\frac{1}{2}) \frac{\Gamma(3/2-n)}{(\frac{1}{2}-n)}} \quad (A-188)$$

Equation A-188 and the original definition in equation A-171 imply that

$$A_n = A_{n-1} = \frac{i(-1)^{n-1} \pi(n-1+\frac{1}{2})}{\Gamma(n-1+3/2)\Gamma(\frac{1}{2}-(n-1))} \quad (A-189)$$

This completes the proof of Lemma A-2.

Previous results were developed on the assumption of a particular form for the expansion of a plane electromagnetic wave. We now verify this expansion. Each component ψ of a plane wave is a solution of the scalar Helmholtz equation

$$\nabla^2 \psi + k^2 \psi = 0. \quad (A-190)$$

In spherical coordinates, equation A-190 takes the form

$$\begin{aligned} \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \frac{\partial \psi}{\partial r}) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta \frac{\partial \psi}{\partial \theta}) \\ + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \psi}{\partial \phi^2} + k^2 \psi = 0. \end{aligned} \quad (A-191)$$

Let us write

$$\psi = R(r)\Theta(\theta)\Phi(\phi), \quad (A-192)$$

$$\psi(r, \theta, \phi) = \sum a_{m,n,p} R_m(r) \Theta_n(\theta) \Phi_p(\phi), \quad (A-193)$$

and observe that $R(r) = R_m(r)$, $\Theta(\theta) = \Theta_n(\theta)$, and $\Phi(\phi) = \Phi_p(\phi)$ imply

$$\begin{aligned} \frac{1}{R(r)} \left[-\frac{1}{r^2} \frac{d}{dr} (r^2 R'(r)) \right] + \frac{1}{\Theta(\theta)} \left[-\frac{1}{r^2 \sin^2 \theta} \frac{d}{d\theta} (\sin^2 \theta \Theta'(\theta)) \right] \\ + \frac{1}{\Phi(\phi)} \left[-\frac{1}{r^2 \sin^2 \theta} \Phi''(\phi) \right] + k^2 = 0. \end{aligned} \quad (A-194)$$

Suppose that electric vector \underline{E} is given by

$$\underline{E} = a_x \exp(-ikz) \exp(i\omega t), \quad (A-195)$$

where \underline{a}_x is the unit base vector directed along the x-axis of a Cartesian coordinate system. The expression for \underline{a}_x in terms of the unit base vectors \underline{e}_r , \underline{e}_θ , and \underline{e}_ϕ in the spherical coordinate system is

$$\underline{a}_x = (\underline{a}_x \cdot \underline{e}_r) \underline{e}_r + (\underline{a}_x \cdot \underline{e}_\theta) \underline{e}_\theta + (\underline{a}_x \cdot \underline{e}_\phi) \underline{e}_\phi \quad (A-196)$$

which, upon replacement of the inner products, becomes

$$\underline{a}_x = \sin\theta \cos\phi \underline{e}_r + \cos\theta \cos\phi \underline{e}_\theta - \sin\phi \underline{e}_\phi. \quad (A-197)$$

Substituting equation A-197 into equation A-195, with z replaced by $r \cos\theta$, yields

$$\begin{aligned} \underline{E} = \sin\theta \cos\phi \exp(-ikr \cos\theta) \exp(i\omega t) \underline{e}_r \\ + \cos\theta \cos\phi \exp(-ikr \cos\theta) \exp(i\omega t) \underline{e}_\theta \\ - \sin\phi \exp(-ikr \cos\theta) \exp(i\omega t) \underline{e}_\phi. \end{aligned} \quad (A-198)$$

Through use of $\underline{E} = \underline{M}_v + i \underline{N}_u$ and equations A-87 through A-90 and A-93 through A-96, the vector \underline{E} can be expressed as

$$\begin{aligned} \underline{E} = & [0\underline{e}_r + \left(\frac{1}{\sin\theta}\right) \frac{\partial v}{\partial \phi} \underline{e}_\theta - \frac{\partial v}{\partial \theta} \underline{e}_\phi] \\ & + i \left\{ \frac{1}{mk} \left[\frac{1}{r} \frac{\partial}{\partial r} (r^2 \frac{\partial u}{\partial r}) + m^2 k^2 r u \right] \underline{e}_r \right. \\ & \left. + \frac{1}{mkr} \frac{\partial^2 (ru)}{\partial r \partial \theta} \underline{e}_\theta + \frac{1}{mk r \sin\theta} \frac{\partial^2 (ru)}{\partial r \partial \phi} \underline{e}_\phi \right\}. \end{aligned} \quad (A-199)$$

Our starting point in finding expressions for u and v is to write down the spherical harmonic expansion of $\exp(-ikr \cos\theta)$. We set

$$\exp(-ikr \cos\theta) = \sum_{n=0}^{\infty} a_n P_n(\cos\theta) j_n(kr), \quad (A-200)$$

where $P_n(\cos\theta)$ and $j_n(kr)$ are the Legendre and spherical Bessel functions respectively. Now $j_n(s)$ has an infinite series expansion

$$j_n(s) = 2^n s^n \sum_{m=0}^{\infty} \frac{(-1)^m (n+m)!}{m! (2n+2m+1)!} s^{2m} \quad (A-201)$$

from which we obtain

$$\left. \frac{d^n}{ds^n} j_n(s) \right|_{s=0} = \frac{2^n (n!)^2}{(2n+1)!}. \quad (A-202)$$

Multiplying both sides of equation A-200 by $P_n(\cos\theta) \sin\theta$, we obtain, upon integrating from 0 to π with respect to θ , the result

$$a_n j_n(kr) \int_0^\pi P_n^2(\cos\theta) \sin\theta d\theta = \int_0^\pi \exp(-ikr\cos\theta) P_n(\cos\theta) \sin\theta d\theta. \quad (\text{A-203})$$

Thus, the relation

$$\int_0^\pi P_n^2(\cos\theta) \sin\theta d\theta = \frac{2}{2n+1} \quad (\text{A-204})$$

implies that

$$a_n j_n(kr) \left(\frac{2}{2n+1}\right) = \int_0^\pi \exp(-ikr\cos\theta) P_n(\cos\theta) \sin\theta d\theta. \quad (\text{A-205})$$

Letting $s = kr$ in equation A-205, we find that

$$a_n j_n(s) \left(\frac{2}{2n+1}\right) = \int_0^\pi \exp(-is\cos\theta) P_n(\cos\theta) \sin\theta d\theta. \quad (\text{A-206})$$

Differentiating both sides of equation A-206 with respect to s n times, setting $s = 0$, and using equation A-202, we obtain

$$a_n \frac{2^n (n!)^2}{(2n+1)!} \left(\frac{2}{2n+1}\right) = \int_0^\pi (-i\cos\theta)^n P_n(\cos\theta) \sin\theta d\theta. \quad (\text{A-207})$$

Consequently,

$$a_n \frac{2^n (n!)^2}{(2n+1)!} \left(\frac{2}{2n+1}\right) = (-i)^n \int_0^\pi \cos^n \theta P_n(\cos\theta) \sin\theta d\theta. \quad (\text{A-208})$$

By equation A-208 and the fact that

$$\int_0^\pi \cos^n \theta P_n(\cos\theta) \sin\theta d\theta = \frac{2^{n+1} (n!)^2}{(2n+1)}, \quad (\text{A-209})$$

we thus obtain

$$a_n = (-1)^n (2n+1). \quad (\text{A-210})$$

Now the functions u and v were previously expanded in terms of products of the associated Legendre functions of the first kind and order one and spherical Bessel functions, $P_n^1(\cos\theta)j_n(kr)$. It thus behooves us to represent $\exp(-ikr\cos\theta)$ in terms of these functions. Differentiation of both sides of equation A-200 with respect to θ results in the relation

$$ikr\sin\theta\exp(-ikr\cos\theta) = - \sum_{n=0}^{\infty} a_n P_n^1(\cos\theta) \sin\theta j_n(kr). \quad (\text{A-211})$$

Using equation A-211 and the fact that

$$P_n^1(x) = \sqrt{1-x^2} P_n^1(x), \quad (\text{A-212})$$

we deduce that

$$-ikr\sin\theta\exp(-ikr\cos\theta) = \sum_{n=0}^{\infty} a_n P_n^1(\cos\theta) j_n(kr). \quad (\text{A-213})$$

Noting that

$$\begin{aligned} r^2 \frac{d^2}{dr^2} j_n(kr) + 2r \frac{d}{dr} j_n(kr) \\ + (k^2 r^2 - n(n+1)) j_n(kr) = 0 \end{aligned} \quad (\text{A-214})$$

and defining u by

$$u = \exp(i\omega t) \cos\phi \sum_{n=1}^{\infty} (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos\theta) \zeta_n'(kr), \quad (\text{A-215})$$

shows that the r component of \underline{E} is, by equations A-198 through A-200, and A-213, given by

$$\begin{aligned} E_r &= \cos\phi \sin\theta \exp(-ikr \cos\theta) \exp(i\omega t) \\ &= \exp(i\omega t) \cos\phi \sum_{n=1}^{\infty} (-i)^n (2n+1) P_n^1(\cos\theta) \left(\frac{j_n(kr)}{-ikr} \right), \end{aligned} \quad (\text{A-216})$$

since by equations A-199 and A-215, we deduce that E_r is given by

$$\begin{aligned} E_r &= \frac{i}{kr} \left[\frac{\partial}{\partial r} (r^2 \frac{\partial u}{\partial r}) + k^2 r^2 u \right] \\ &= i \exp(i\omega t) \cos\phi \sum_{n=1}^{\infty} (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos\theta) \frac{n(n+1)}{kr} j_n(kr) \\ &= \exp(i\omega t) \cos\phi \sum_{n=1}^{\infty} (-i)^n (2n+1) P_n^1(\cos\theta) \left(\frac{j_n(kr)}{-ikr} \right) \end{aligned} \quad (\text{A-217})$$

which is exactly the right side of equation A-216.

Now an analysis of the coefficient of \underline{e}_θ in equation A-198 is in order. An essential step toward this goal is to expand the function

$$F(r, \theta) = \exp(-ikr \cos\theta) \cos\theta. \quad (\text{A-218})$$

Differentiating both sides of equation A-200 with respect to r , we find that

$$-i \cos\theta \exp(-ikr \cos\theta) = \sum_{n=0}^{\infty} a_n P_n^1(\cos\theta) j_n'(kr). \quad (\text{A-219})$$

Since $\underline{E} = \underline{M}_v + i \underline{N}_u$ and setting $m = 1$, the coefficient of \underline{e}_θ equation A-199 is given by

$$\frac{1}{\sin\theta} \frac{\partial v}{\partial \phi} + i \left(\frac{1}{k} \frac{\partial^2 u}{\partial r \partial \theta} + \frac{1}{kr} \frac{\partial u}{\partial \theta} \right) . \quad (\text{A-220})$$

One might be tempted to set

$$\frac{\partial v}{\partial \phi} = u \quad (\text{A-221})$$

and use equation A-215 to deduce that

$$v = \exp(i\omega t) \sin\phi \sum_{n=1}^{\infty} (-i)^n \frac{2n+1}{n(n+1)} P_n^1(\cos\theta) j_n(kr) . \quad (\text{A-222})$$

Also one might attempt to prove that

$$\begin{aligned} \frac{1}{\sin\theta} \frac{\partial v}{\partial \phi} + i \left(\frac{1}{k} \frac{\partial^2 u}{\partial r \partial \theta} + \frac{1}{kr} \frac{\partial u}{\partial \theta} \right) &= \exp(i\omega t) \sin\theta \cos\phi \\ &\cdot \exp(-ikr \cos\theta) \end{aligned} \quad (\text{A-223})$$

by substitution of series. However, there is a much easier method.

This involves seeking an expression for u in the form

$$u = \frac{B(\theta, \phi) \exp(-ikr \cos\theta)}{r} . \quad (\text{A-224})$$

Then

$$\frac{\partial u}{\partial r} = B(\theta, \phi) \exp(-ikr \cos\theta) \left(-\frac{1}{r^2} - \frac{ik \cos\theta}{r} \right) , \quad (\text{A-225})$$

$$r^2 \frac{\partial u}{\partial r} = B(\theta, \phi) \exp(-ikr \cos\theta) (-1 - ikr \cos\theta) , \quad (\text{A-226})$$

$$\begin{aligned} \frac{\partial}{\partial r} \left(r^2 \frac{\partial u}{\partial r} \right) &= B(\theta, \phi) \exp(-ikr \cos\theta) [(-ik \cos\theta)(-1 - ikr \cos\theta) \\ &\quad - ik \cos\theta] . \end{aligned} \quad (\text{A-227})$$

Equations A-217, A-224, and A-227 lead to

$$\begin{aligned}
 E_r &= \frac{i}{kr} \left[\frac{\partial}{\partial r} (r^2 \frac{\partial u}{\partial r}) + k^2 r^2 u \right] \\
 &= B(\theta, \phi) \exp(-ikr \cos \theta) \cdot \left\{ \frac{\cos \theta}{r} [(-1 - ikr \cos \theta) + 1] \right\} \\
 &\quad + ikr [B(\theta, \phi) \exp(-ikr \cos \theta) / r] \\
 &= ik(1 - \cos^2 \theta) B(\theta, \phi) \exp(-ikr \cos \theta) \\
 &= ik \sin^2 \theta B(\theta, \phi) \exp(-ikr \cos \theta). \quad (A-228)
 \end{aligned}$$

Now we desire to find $B(\theta, \phi)$ so that equation A-228 becomes

$$\frac{i}{kr} \left[\frac{\partial}{\partial r} (r^2 \frac{\partial u}{\partial r}) + k^2 r^2 u \right] = \exp(i\omega t) \sin \theta \cos \phi \exp(-kr \cos \theta). \quad (A-229)$$

The choice

$$B(\theta, \phi) = \frac{\cos \phi}{ik \sin \theta} \exp(i\omega t) \quad (A-230)$$

will serve our purpose. Thus, we must attempt to expand

$$u = \exp(i\omega t) \cos \phi \frac{\exp(-ikr \cos \theta)}{ikr \sin \theta}. \quad (A-231)$$

The method of undetermined coefficients will enable us to obtain equation A-215 by using equation A-229.

Let

$$v = \exp(i\omega t) \sin \phi \frac{\exp(-ikr \cos \theta)}{ikr \sin \theta}, \quad (A-232)$$

Let $a(t) = \exp(i\omega t)$, and form the derivatives

$$\frac{\partial v}{\partial \phi} = a(t) \exp(-ikr \cos \theta) \frac{\cos \phi}{ikr \sin \theta}, \quad (\text{A-233})$$

$$\frac{\partial u}{\partial \theta} = a(t) \exp(-ikr \cos \theta) \cos \phi \left(\frac{-\cos \theta}{ikr \sin^2 \theta} + 1 \right), \quad (\text{A-234})$$

$$\begin{aligned} \frac{\partial^2 u}{\partial r \partial \theta} = a(t) \exp(-ikr \cos \theta) \cos \phi & \left(\frac{ik \cos^2 \theta}{ir \sin^2 \theta} - ik \cos \theta \right. \\ & \left. + \frac{\cos \theta}{ikr^2 \sin^2 \theta} \right). \end{aligned} \quad (\text{A-235})$$

Construct the left member of equation A-223 by using equations A-233 through A-235 and find that

$$\begin{aligned} & \frac{1}{\sin \theta} \frac{\partial v}{\partial \phi} + i \left(\frac{1}{k} \frac{\partial^2 u}{\partial r \partial \theta} + \frac{1}{kr} \frac{\partial u}{\partial \theta} \right) \\ &= a(t) \exp(-ikr \cos \theta) \cos \phi \left(\frac{1}{ikr \sin^2 \theta} + \frac{icos^2 \theta}{kr \sin^2 \theta} \right. \\ & \quad \left. + \cos \theta + \frac{\cos \theta}{k^2 r^2 \sin^2 \theta} - \frac{\cos \theta}{k^2 r^2 \sin^2 \theta} + \frac{1}{kr} \right) \\ &= a(t) \exp(-ikr \cos \theta) \cos \phi \left(\frac{1}{ikr \sin^2 \theta} + \frac{icos^2 \theta}{kr \sin^2 \theta} \right. \\ & \quad \left. + \frac{isin^2 \theta}{kr \sin^2 \theta} + \cos \theta \right) \\ &= a(t) \cos \theta \cos \phi \exp(-ikr \cos \theta) = E_{\theta}. \end{aligned} \quad (\text{A-236})$$

This establishes the validity of equation A-223. Finally, to complete the argument, we must show that

$$-\frac{\partial v}{\partial \phi} + \frac{i}{k r \sin \theta} \frac{\partial^2}{\partial r \partial \phi} (ru) = -\exp(i\omega t) \sin \phi \exp(-ikr \cos \theta), \quad (\text{A-237})$$

We have from equation A-232 that

$$\frac{\partial v}{\partial \theta} = \exp(i\omega t) \sin \phi \exp(-ikr \cos \theta) \left(\frac{-\cos \theta}{ik r \sin^2 \theta} + 1 \right). \quad (\text{A-238})$$

Also by equation A-231,

$$\frac{\partial^2}{\partial r \partial \phi} (ru) = \exp(i\omega t) \sin \phi \exp(-ikr \cos \theta) \left(\frac{\cos \theta}{\sin \theta} \right). \quad (\text{A-239})$$

Thus, equation A-239 implies that

$$\frac{1}{k r \sin \theta} \frac{\partial^2}{\partial r \partial \phi} (ru) = \exp(i\omega t) \sin \phi \exp(-ikr \cos \theta) \left(\frac{\cos \theta}{k r \sin^2 \theta} \right). \quad (\text{A-240})$$

Consequently, combining equations A-238 and A-240, we deduce that

$$\begin{aligned} -\frac{\partial v}{\partial \theta} + \frac{i}{k r \sin \theta} \frac{\partial^2}{\partial r \partial \phi} (ru) &= -\exp(i\omega t) \sin \phi \exp(-ikr \cos \theta) \\ &\quad \cdot \left(1 + \frac{\cos \theta}{ik r \sin^2 \theta} + \frac{i \cos \theta}{k r \sin^2 \theta} \right) \\ &= -\exp(i\omega t) \sin \phi \exp(-ikr \cos \theta) = E_{\phi}. \quad (\text{A-241}) \end{aligned}$$

This completes the derivation. The following lemma has been proved:

Lemma A-3. If u and v are defined by equations A-231 and A-232, respectively, then u satisfies equation A-215, v satisfies equation A-222, and u and v satisfy equation A-199, where vector \underline{E} is defined by equation A-198.

In proving that the u defined by equation A-231 has the representation A-215, we used some properties of Legendre polynomials which are consequences of the following lemma:

Lemma A-4. For every positive integer n ,

$$\int_0^\pi \cos^n \theta P_n(\cos \theta) \sin \theta d\theta = \frac{2^{n+1} (n!)^2}{(2n+1)!} \quad (\text{A-242})$$

and

$$\int_0^\pi [P_n(\cos \theta)]^2 \sin \theta d\theta = \frac{2}{2n+1} \quad (\text{A-243})$$

Proof of Lemma A-4. We use Rodrigues' formula

$$P_n(x) = \frac{1}{2^n n!} \frac{d^n}{dx^n} (x^2 - 1)^n \quad (\text{A-244})$$

and observe that the substitution $x = \cos \theta$ implies that

$$\begin{aligned} \int_0^\pi \cos^n \theta P_n(\cos \theta) \sin \theta d\theta &= \frac{1}{2^n n!} \int_{-1}^1 x^n \frac{d^n}{dx^n} (x^2 - 1)^n dx \\ \int_0^\pi \cos^n \theta P_n(\cos \theta) \sin \theta d\theta &= \frac{(-1)^n}{2^n n!} \int_{-1}^1 \left(\frac{d^n}{dx^n} x^n \right) (x^2 - 1)^n dx \\ &= \frac{(-1)^n}{2^n} \int_{-1}^1 (x-1)^n (x+1)^n dx \quad (\text{A-245}) \end{aligned}$$

Continuing the integration by parts of the right side of equation A-245, we conclude that

$$\begin{aligned}
 \int_0^\pi \cos^n \theta P_n(\cos \theta) \sin \theta d\theta &= \frac{(-1)^{2n}}{2^n (n+1)(n+2) \dots (n+n)} \int_{-1}^1 \left[\left(\frac{d^n}{dx^n} \right) (x-1)^n \right] (x+1)^{2n} dx \\
 &= \frac{2^{2n+1} (n!)^2}{2^n (2n+1)!} \\
 &= \frac{2^{n+1} (n!)^2}{(2n+1)!} . \quad (A-246)
 \end{aligned}$$

To complete the proof of the lemma, observe that

$$\begin{aligned}
 \int_0^\pi [P_n(\cos \theta)]^2 \sin \theta d\theta &= \int_{-1}^1 [P_n(x)]^2 dx \\
 &= \frac{1}{2^{2n} (n!)^2} \int_{-1}^1 \left[\frac{d^n}{dx^n} (x^2-1)^n \right] \left[\frac{d^n}{dx^n} (x^2-1)^n \right] dx \\
 &= \frac{(-1)^n}{2^{2n} (n!)^2} \int_{-1}^1 \left[\frac{d^{2n}}{dx^{2n}} (x^2-1)^n \right] (x^2-1)^n dx. \quad (A-247)
 \end{aligned}$$

Since

$$\frac{d^{2n}}{dx^{2n}} (x^2-1)^n = (2n)! \quad (A-248)$$

equation A-245 implies that

$$\int_0^{\pi} [P_n(\cos\theta)]^2 \sin\theta d\theta = \frac{(-1)^n (2n)!}{2^{2n} (n!)^2} \int_{-1}^1 (x^2-1)^n dx. \quad (\text{A-249})$$

But another integration by parts shows that

$$\begin{aligned} \int_{-1}^1 (x^2-1)^n dx &= \int_{-1}^1 (x-1)^n (x+1)^n dx \\ &= \frac{2^{n+1} (n!)^2}{(2n+1)!} . \end{aligned} \quad (\text{A-250})$$

Substituting equation A-250 into equation A-249, we obtain equation A-243.

APPENDIX B

SAMPLE PROBLEMS WITH COMPUTER RESULTS

SAMPLE PROBLEM 1 DECK SETUP

CARD 1 (CONTROL PARAMETERS)

1.000E3 5.986E1 1.00152E0 1.000E0 1.000E1 0.000E0 66 1 8 8 8

CARDS 2 - (NOC+1) (DATA CARDS)

8.660E-01 5.474E+01 4.500E+01
1.658E+00 2.524E+01 4.500E+01
2.598E+00 1.579E+01 4.500E+01
3.571E+00 1.142E+01 4.500E+01
4.555E+00 8.930E+00 4.500E+01
1.658E+00 7.245E+01 7.157E+01
2.179E+00 4.651E+01 7.157E+01
2.958E+00 3.231E+01 7.157E+01
3.841E+00 2.431E+01 7.157E+01
4.770E+00 1.936E+01 7.157E+01
2.598E+00 7.890E+01 7.869E+01
2.958E+00 5.953E+01 7.869E+01
3.571E+00 4.556E+01 7.869E+01
4.330E+00 3.607E+01 7.869E+01
3.571E+00 8.195E+01 8.187E+01
3.841E+00 6.701E+01 8.187E+01
4.330E+00 5.474E+01 8.187E+01
4.975E+00 4.529E+01 8.187E+01
4.555E+00 8.370E+01 8.366E+01
4.770E+00 7.167E+01 8.366E+01
1.658E+00 7.245E+01 1.843E+01
2.179E+00 4.651E+01 1.843E+01
2.958E+00 3.231E+01 1.843E+01
3.841E+00 2.431E+01 1.843E+01
4.770E+00 1.936E+01 1.843E+01
2.179E+00 7.674E+01 4.500E+01
2.598E+00 5.474E+01 4.500E+01
3.279E+00 4.032E+01 4.500E+01
4.093E+00 3.122E+01 4.500E+01
4.975E+00 2.524E+01 4.500E+01
2.958E+00 8.027E+01 5.904E+01
3.279E+00 6.277E+01 5.904E+01
3.841E+00 4.939E+01 5.904E+01
4.555E+00 3.979E+01 5.904E+01
3.841E+00 8.252E+01 6.680E+01
4.093E+00 6.850E+01 6.680E+01
4.555E+00 5.671E+01 6.680E+01
4.770E+00 8.398E+01 7.157E+01
4.975E+00 7.245E+01 7.157E+01
2.598E+00 7.890E+01 1.131E+01
2.958E+00 5.953E+01 1.131E+01
3.571E+00 4.556E+01 1.131E+01
4.330E+00 3.607E+01 1.131E+01
2.958E+00 8.027E+01 3.096E+01
3.279E+00 6.277E+01 3.096E+01
3.841E+00 4.939E+01 3.096E+01
4.555E+00 3.979E+01 3.096E+01
3.571E+00 8.195E+01 4.500E+01
3.841E+00 6.701E+01 4.500E+01
4.330E+00 5.474E+01 4.500E+01
4.330E+00 8.337E+01 5.446E+01
4.555E+00 7.077E+01 5.446E+01
3.571E+00 8.195E+01 8.130E+00
3.841E+00 6.701E+01 8.130E+00
4.330E+00 5.474E+01 8.130E+00
4.975E+00 4.529E+01 8.130E+00
3.841E+00 8.252E+01 2.320E+01
4.093E+00 6.850E+01 2.320E+01
4.555E+00 5.671E+01 2.320E+01
4.330E+00 6.337E+01 3.554E+01
4.555E+00 7.077E+01 3.554E+01
4.975E+00 8.423E+01 4.500E+01
4.555E+00 8.370E+01 6.340E+00
4.770E+00 7.167E+01 6.340E+00
4.770E+00 8.398E+01 1.843E+01
4.975E+00 7.245E+01 1.843E+01

TERMINATION CARD

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SAMPLE PROBLEM 1

DEPOSITION OF POWER INSIDE A HOMOGENEOUS SPHERE IMMERSED IN AN ELECTROMAGNETIC FIELD

FREQUENCY = 1000.00 MHZ WAVELENGTH = 29.97925 CM FIELD STRENGTH = 1.00 V/H
 CONDUCTIVITY = 1.001520 MHO/CM RELATIVE DIELECTRIC CONSTANT = 59.86 DIAMETER = 10.00 CM TIME = 0.0 SEC
 REFRACTIVE INDEX = 7.92226D+00 -1.15231D+00

THE HIGHEST ORDER NEUMANN FUNCTION COMPUTED IS OF ORDER = 21 SIZE PARAMETER = 1.04792D+00

RATIO OF THE FIRST CALCULATION OF THE ZERO ORDER SPHERICAL BESSEL FUNCTION OF ARGUMENT X
 TO THE VALUE OF $\sin(X)/X = 1.00000D+00$

RATIO OF THE FIRST CALCULATION OF THE ZERO ORDER SPHERICAL BESSEL FUNCTION OF ARGUMENT Z
 TO THE VALUE OF $\sin(Z)/Z = 1.00000D+00$ 7.13511D-16

INTERNAL POINT:	RADIUS =	0.865 CM	THETA =	54.74 DEG	PHI =	45.00 DEG	ABSORBED POWER DENSITY =	0.09375857 WATTS/CM ³
INTERNAL POINT:	RADIUS =	1.658 CM	THETA =	25.24 DEG	PHI =	45.00 DEG	ABSORBED POWER DENSITY =	0.03009165 WATTS/CM ³
INTERNAL POINT:	RADIUS =	2.598 CM	THETA =	15.79 DEG	PHI =	45.00 DEG	ABSORBED POWER DENSITY =	0.05211915 WATTS/CM ³
INTERNAL POINT:	RADIUS =	3.571 CM	THETA =	11.42 DEG	PHI =	45.00 DEG	ABSORBED POWER DENSITY =	0.00718193 WATTS/CM ³
INTERNAL POINT:	RADIUS =	4.555 CM	THETA =	8.93 DEG	PHI =	45.00 DEG	ABSORBED POWER DENSITY =	0.02350504 WATTS/CM ³
INTERNAL POINT:	RADIUS =	1.658 CM	THETA =	72.45 DEG	PHI =	71.57 DEG	ABSORBED POWER DENSITY =	0.01844696 WATTS/CM ³
INTERNAL POINT:	RADIUS =	2.179 CM	THETA =	46.51 DEG	PHI =	71.57 DEG	ABSORBED POWER DENSITY =	0.01423773 WATTS/CM ³
INTERNAL POINT:	RADIUS =	2.958 CM	THETA =	32.31 DEG	PHI =	71.57 DEG	ABSORBED POWER DENSITY =	0.01821150 WATTS/CM ³
INTERNAL POINT:	RADIUS =	3.841 CM	THETA =	24.31 DEG	PHI =	71.57 DEG	ABSORBED POWER DENSITY =	0.00658279 WATTS/CM ³
INTERNAL POINT:	RADIUS =	4.770 CM	THETA =	19.36 DEG	PHI =	71.57 DEG	ABSORBED POWER DENSITY =	0.01929612 WATTS/CM ³
INTERNAL POINT:	RADIUS =	2.598 CM	THETA =	78.90 DEG	PHI =	78.93 DEG	ABSORBED POWER DENSITY =	0.01421962 WATTS/CM ³
INTERNAL POINT:	RADIUS =	2.958 CM	THETA =	59.53 DEG	PHI =	78.93 DEG	ABSORBED POWER DENSITY =	0.00863334 WATTS/CM ³
INTERNAL POINT:	RADIUS =	3.571 CM	THETA =	45.56 DEG	PHI =	78.93 DEG	ABSORBED POWER DENSITY =	0.00203956 WATTS/CM ³
INTERNAL POINT:	RADIUS =	4.330 CM	THETA =	36.07 DEG	PHI =	78.93 DEG	ABSORBED POWER DENSITY =	0.01071324 WATTS/CM ³
INTERNAL POINT:	RADIUS =	3.571 CM	THETA =	81.95 DEG	PHI =	81.97 DEG	ABSORBED POWER DENSITY =	0.00608587 WATTS/CM ³
INTERNAL POINT:	RADIUS =	3.841 CM	THETA =	67.01 DEG	PHI =	81.97 DEG	ABSORBED POWER DENSITY =	0.00286658 WATTS/CM ³
INTERNAL POINT:	RADIUS =	4.330 CM	THETA =	58.74 DEG	PHI =	81.97 DEG	ABSORBED POWER DENSITY =	0.00525896 WATTS/CM ³
INTERNAL POINT:	RADIUS =	4.975 CM	THETA =	45.29 DEG	PHI =	81.97 DEG	ABSORBED POWER DENSITY =	0.00696834 WATTS/CM ³
INTERNAL POINT:	RADIUS =	4.555 CM	THETA =	83.70 DEG	PHI =	83.66 DEG	ABSORBED POWER DENSITY =	0.00902432 WATTS/CM ³
INTERNAL POINT:	RADIUS =	4.770 CM	THETA =	71.67 DEG	PHI =	83.66 DEG	ABSORBED POWER DENSITY =	0.00847169 WATTS/CM ³
INTERNAL POINT:	RADIUS =	1.658 CM	THETA =	72.45 DEG	PHI =	18.43 DEG	ABSORBED POWER DENSITY =	0.07920915 WATTS/CM ³
INTERNAL POINT:	RADIUS =	2.179 CM	THETA =	46.51 DEG	PHI =	18.43 DEG	ABSORBED POWER DENSITY =	0.05024135 WATTS/CM ³
INTERNAL POINT:	RADIUS =	2.958 CM	THETA =	32.31 DEG	PHI =	18.43 DEG	ABSORBED POWER DENSITY =	0.02015898 WATTS/CM ³
INTERNAL POINT:	RADIUS =	3.841 CM	THETA =	24.31 DEG	PHI =	18.43 DEG	ABSORBED POWER DENSITY =	0.00823060 WATTS/CM ³
INTERNAL POINT:	RADIUS =	4.770 CM	THETA =	19.36 DEG	PHI =	18.43 DEG	ABSORBED POWER DENSITY =	0.02027197 WATTS/CM ³
INTERNAL POINT:	RADIUS =	2.179 CM	THETA =	76.74 DEG	PHI =	45.00 DEG	ABSORBED POWER DENSITY =	0.01373777 WATTS/CM ³
INTERNAL POINT:	RADIUS =	2.598 CM	THETA =	54.74 DEG	PHI =	45.00 DEG	ABSORBED POWER DENSITY =	0.01572690 WATTS/CM ³
INTERNAL POINT:	RADIUS =	3.279 CM	THETA =	40.32 DEG	PHI =	45.00 DEG	ABSORBED POWER DENSITY =	0.00521308 WATTS/CM ³
INTERNAL POINT:	RADIUS =	4.093 CM	THETA =	31.22 DEG	PHI =	45.00 DEG	ABSORBED POWER DENSITY =	0.01125208 WATTS/CM ³
INTERNAL POINT:	RADIUS =	4.975 CM	THETA =	25.24 DEG	PHI =	45.00 DEG	ABSORBED POWER DENSITY =	0.01501013 WATTS/CM ³
INTERNAL POINT:	RADIUS =	2.958 CM	THETA =	80.27 DEG	PHI =	59.04 DEG	ABSORBED POWER DENSITY =	0.01176939 WATTS/CM ³
INTERNAL POINT:	RADIUS =	3.279 CM	THETA =	62.77 DEG	PHI =	59.04 DEG	ABSORBED POWER DENSITY =	0.00532103 WATTS/CM ³
INTERNAL POINT:	RADIUS =	3.841 CM	THETA =	49.39 DEG	PHI =	59.04 DEG	ABSORBED POWER DENSITY =	0.00428390 WATTS/CM ³
INTERNAL POINT:	RADIUS =	4.555 CM	THETA =	39.79 DEG	PHI =	59.04 DEG	ABSORBED POWER DENSITY =	0.01136360 WATTS/CM ³
INTERNAL POINT:	RADIUS =	3.841 CM	THETA =	82.52 DEG	PHI =	66.00 DEG	ABSORBED POWER DENSITY =	0.00819845 WATTS/CM ³
INTERNAL POINT:	RADIUS =	4.093 CM	THETA =	68.50 DEG	PHI =	66.00 DEG	ABSORBED POWER DENSITY =	0.00535788 WATTS/CM ³
INTERNAL POINT:	RADIUS =	4.555 CM	THETA =	56.71 DEG	PHI =	66.00 DEG	ABSORBED POWER DENSITY =	0.00632622 WATTS/CM ³
INTERNAL POINT:	RADIUS =	4.770 CM	THETA =	83.98 DEG	PHI =	71.57 DEG	ABSORBED POWER DENSITY =	0.00936273 WATTS/CM ³
INTERNAL POINT:	RADIUS =	4.975 CM	THETA =	72.45 DEG	PHI =	71.57 DEG	ABSORBED POWER DENSITY =	0.00625014 WATTS/CM ³
INTERNAL POINT:	RADIUS =	2.598 CM	THETA =	78.90 DEG	PHI =	11.31 DEG	ABSORBED POWER DENSITY =	0.00915834 WATTS/CM ³
INTERNAL POINT:	RADIUS =	2.958 CM	THETA =	59.53 DEG	PHI =	11.31 DEG	ABSORBED POWER DENSITY =	0.00502638 WATTS/CM ³
INTERNAL POINT:	RADIUS =	3.571 CM	THETA =	45.56 DEG	PHI =	11.31 DEG	ABSORBED POWER DENSITY =	0.00522500 WATTS/CM ³
INTERNAL POINT:	RADIUS =	4.330 CM	THETA =	36.07 DEG	PHI =	11.31 DEG	ABSORBED POWER DENSITY =	0.01555046 WATTS/CM ³

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INTERNAL POINT:	RADIUS =	2.958 CM	THETA =	80.27 DEG	PHI =	30.96 DEG	ABSORBED POWER DENSITY =	0.01024710 WATTS/M**3
INTERNAL POINT:	RADIUS =	3.279 CM	THETA =	62.77 DEG	PHI =	30.96 DEG	ABSORBED POWER DENSITY =	0.00568853 WATTS/M**3
INTERNAL POINT:	RADIUS =	3.641 CM	THETA =	49.39 DEG	PHI =	30.96 DEG	ABSORBED POWER DENSITY =	0.00808898 WATTS/M**3
INTERNAL POINT:	RADIUS =	4.555 CM	THETA =	39.79 DEG	PHI =	30.96 DEG	ABSORBED POWER DENSITY =	0.01323507 WATTS/M**3
INTERNAL POINT:	RADIUS =	3.571 CM	THETA =	81.95 DEG	PHI =	45.00 DEG	ABSORBED POWER DENSITY =	0.01331812 WATTS/M**3
INTERNAL POINT:	RADIUS =	3.841 CM	THETA =	67.01 DEG	PHI =	45.00 DEG	ABSORBED POWER DENSITY =	0.00864617 WATTS/M**3
INTERNAL POINT:	RADIUS =	4.330 CM	THETA =	54.74 DEG	PHI =	45.00 DEG	ABSORBED POWER DENSITY =	0.00852068 WATTS/M**3
INTERNAL POINT:	RADIUS =	4.330 CM	THETA =	83.37 DEG	PHI =	54.46 DEG	ABSORBED POWER DENSITY =	0.00978453 WATTS/M**3
INTERNAL POINT:	RADIUS =	4.555 CM	THETA =	70.77 DEG	PHI =	54.46 DEG	ABSORBED POWER DENSITY =	0.00690478 WATTS/M**3
INTERNAL POINT:	RADIUS =	3.571 CM	THETA =	81.95 DEG	PHI =	8.13 DEG	ABSORBED POWER DENSITY =	0.01976837 WATTS/M**3
INTERNAL POINT:	RADIUS =	3.841 CM	THETA =	67.01 DEG	PHI =	8.13 DEG	ABSORBED POWER DENSITY =	0.01442377 WATTS/M**3
INTERNAL POINT:	RADIUS =	4.330 CM	THETA =	54.74 DEG	PHI =	8.13 DEG	ABSORBED POWER DENSITY =	0.01178249 WATTS/M**3
INTERNAL POINT:	RADIUS =	4.330 CM	THETA =	83.37 DEG	PHI =	8.13 DEG	ABSORBED POWER DENSITY =	0.00748877 WATTS/M**3
INTERNAL POINT:	RADIUS =	4.555 CM	THETA =	70.77 DEG	PHI =	23.20 DEG	ABSORBED POWER DENSITY =	0.01738899 WATTS/M**3
INTERNAL POINT:	RADIUS =	4.555 CM	THETA =	69.50 DEG	PHI =	23.20 DEG	ABSORBED POWER DENSITY =	0.01299063 WATTS/M**3
INTERNAL POINT:	RADIUS =	4.330 CM	THETA =	56.74 DEG	PHI =	23.20 DEG	ABSORBED POWER DENSITY =	0.00864552 WATTS/M**3
INTERNAL POINT:	RADIUS =	4.555 CM	THETA =	83.37 DEG	PHI =	35.54 DEG	ABSORBED POWER DENSITY =	0.01132341 WATTS/M**3
INTERNAL POINT:	RADIUS =	4.975 CM	THETA =	70.77 DEG	PHI =	35.54 DEG	ABSORBED POWER DENSITY =	0.00757644 WATTS/M**3
INTERNAL POINT:	RADIUS =	4.555 CM	THETA =	86.23 DEG	PHI =	45.00 DEG	ABSORBED POWER DENSITY =	0.01062630 WATTS/M**3
INTERNAL POINT:	RADIUS =	4.770 CM	THETA =	83.70 DEG	PHI =	6.34 DEG	ABSORBED POWER DENSITY =	0.01101610 WATTS/M**3
INTERNAL POINT:	RADIUS =	4.770 CM	THETA =	71.67 DEG	PHI =	6.34 DEG	ABSORBED POWER DENSITY =	0.00690943 WATTS/M**3
INTERNAL POINT:	RADIUS =	4.770 CM	THETA =	83.98 DEG	PHI =	18.43 DEG	ABSORBED POWER DENSITY =	0.01083597 WATTS/M**3
INTERNAL POINT:	RADIUS =	4.975 CM	THETA =	72.45 DEG	PHI =	18.43 DEG	ABSORBED POWER DENSITY =	0.00709315 WATTS/M**3

TOTAL ABSORBED POWER= 1.08523D-05 WATTS

AVERAGE ABSORBED POWER DENSITY= 2.07264D-02 WATTS/M**3

APPROXIMATE EXECUTION TIME = 0.17 CPU MINUTE

SAMPLE PROBLEM 2 DECK SETUP

CARD 1 (CONTROL PARAMETERS)

1.000E3 5.986E1 1.00152E0 1.000E0 1.000E1 0.000E0 101 0 0 0 0

CARDS 2 - (NOC+1) (DATA CARDS)

0.1E-3	0.0E0	0.0E0
0.1E0	0.0E0	0.0E0
0.2E0	0.0E0	0.0E0
0.3E0	0.0E0	0.0E0
0.4E0	0.0E0	0.0E0
0.5E0	0.0E0	0.0E0
0.6E0	0.0E0	0.0E0
0.7E0	0.0E0	0.0E0
0.8E0	0.0E0	0.0E0
0.9E0	0.0E0	0.0E0
1.0E0	0.0E0	0.0E0
1.1E0	0.0E0	0.0E0
1.2E0	0.0E0	0.0E0
1.3E0	0.0E0	0.0E0
1.4E0	0.0E0	0.0E0
1.5E0	0.0E0	0.0E0
1.6E0	0.0E0	0.0E0
1.7E0	0.0E0	0.0E0
1.8E0	0.0E0	0.0E0
1.9E0	0.0E0	0.0E0
2.0E0	0.0E0	0.0E0
2.1E0	0.0E0	0.0E0
2.2E0	0.0E0	0.0E0
2.3E0	0.0E0	0.0E0
2.4E0	0.0E0	0.0E0
2.5E0	0.0E0	0.0E0
2.6E0	0.0E0	0.0E0
2.7E0	0.0E0	0.0E0
2.8E0	0.0E0	0.0E0
2.9E0	0.0E0	0.0E0
3.0E0	0.0E0	0.0E0
3.1E0	0.0E0	0.0E0
3.2E0	0.0E0	0.0E0
3.3E0	0.0E0	0.0E0
3.4E0	0.0E0	0.0E0
3.5E0	0.0E0	0.0E0
3.6E0	0.0E0	0.0E0
3.7E0	0.0E0	0.0E0
3.8E0	0.0E0	0.0E0
3.9E0	0.0E0	0.0E0
4.0E0	0.0E0	0.0E0
4.1E0	0.0E0	0.0E0
4.2E0	0.0E0	0.0E0
4.3E0	0.0E0	0.0E0
4.4E0	0.0E0	0.0E0
4.5E0	0.0E0	0.0E0
4.6E0	0.0E0	0.0E0
4.7E0	0.0E0	0.0E0
4.8E0	0.0E0	0.0E0
4.9E0	0.0E0	0.0E0

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5.0E0	0.0E0	0.0E0
0.1E0	180.0E0	0.0E0
0.2E0	180.0E0	0.0E0
0.3E0	180.0E0	0.0E0
0.4E0	180.0E0	0.0E0
0.5E0	180.0E0	0.0E0
0.6E0	180.0E0	0.0E0
0.7E0	180.0E0	0.0E0
0.8E0	180.0E0	0.0E0
0.9E0	180.0E0	0.0E0
1.0E0	180.0E0	0.0E0
1.1E0	180.0E0	0.0E0
1.2E0	180.0E0	0.0E0
1.3E0	180.0E0	0.0E0
1.4E0	180.0E0	0.0E0
1.5E0	180.0E0	0.0E0
1.6E0	180.0E0	0.0E0
1.7E0	180.0E0	0.0E0
1.8E0	180.0E0	0.0E0
1.9E0	180.0E0	0.0E0
2.0E0	180.0E0	0.0E0
2.1E0	180.0E0	0.0E0
2.2E0	180.0E0	0.0E0
2.3E0	180.0E0	0.0E0
2.4E0	180.0E0	0.0E0
2.5E0	180.0E0	0.0E0
2.6E0	180.0E0	0.0E0
2.7E0	180.0E0	0.0E0
2.8E0	180.0E0	0.0E0
2.9E0	180.0E0	0.0E0
3.0E0	180.0E0	0.0E0
3.1E0	180.0E0	0.0E0
3.2E0	180.0E0	0.0E0
3.3E0	180.0E0	0.0E0
3.4E0	180.0E0	0.0E0
3.5E0	180.0E0	0.0E0
3.6E0	180.0E0	0.0E0
3.7E0	180.0E0	0.0E0
3.8E0	180.0E0	0.0E0
3.9E0	180.0E0	0.0E0
4.0E0	180.0E0	0.0E0
4.1E0	180.0E0	0.0E0
4.2E0	180.0E0	0.0E0
4.3E0	180.0E0	0.0E0
4.4E0	180.0E0	0.0E0
4.5E0	180.0E0	0.0E0
4.6E0	180.0E0	0.0E0
4.7E0	180.0E0	0.0E0
4.8E0	180.0E0	0.0E0
4.9E0	180.0E0	0.0E0
5.0E0	180.0E0	0.0E0

TERMINATION CARD

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CASE 2 PROBLEM 2

DEPOSITION OF POWER INSIDE A HOMOGENEOUS SPHERE IMMERSED IN AN ELECTROMAGNETIC FIELD

FREQUENCY = 1000.00 MHZ WAVELENGTH = 29.97925 CM FIELD STRENGTH = 1.00 V/M
 CONDUCTIVITY = 1.001520 MHO/M RELATIVE DIELECTRIC CONSTANT = 59.86 DIAMETER = 10.00 CM TIME = 0.0 SEC

REACTIVE INDEX = 7.822260+00 -1.152310+00

THE HIGHEST ORDER NEUMANN FUNCTION COMPUTED IS OF ORDER = 21 SIZE PARAMETER = 1.007920+00

RATIO OF THE FIRST CALCULATION OF THE ZERO ORDER SPHERICAL BESSEL FUNCTION OF ARGUMENT X
 TO THE VALUE OF $\sin(x)/x = 1.000000+00$

RATIO OF THE FIRST CALCULATION OF THE ZERO ORDER SPHERICAL BESSEL FUNCTION OF ARGUMENT X
 TO THE VALUE OF $\sin(x)/x = 1.000000+00 7.135110-15$

INTERNAL POINT:	RADIUS =	0.0001CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	0.100 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	0.200 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	0.300 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	0.400 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	0.500 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	0.600 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	0.700 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	0.800 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	0.900 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	1.000 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	1.100 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	1.200 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	1.300 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	1.400 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	1.500 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	1.600 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	1.700 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	1.800 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	1.900 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	2.000 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	2.100 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	2.200 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	2.300 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	2.400 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	2.500 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	2.600 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	2.700 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	2.800 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	2.900 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	3.000 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	3.100 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	3.200 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	3.300 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	3.400 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	3.500 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	3.600 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	3.700 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	3.800 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	3.900 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	4.000 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	4.100 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03
INTERNAL POINT:	RADIUS =	4.200 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ABSORBED POWER DENSITY =	0.1692820E-03

INTERVAL POINT:	RADIUS =	4.300 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.02042313	WATTS/CM ²
INTERVAL POINT:	RADIUS =	4.400 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.02232636	WATTS/CM ²
INTERVAL POINT:	RADIUS =	4.500 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.02368835	WATTS/CM ²
INTERVAL POINT:	RADIUS =	4.600 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.02431844	WATTS/CM ²
INTERVAL POINT:	RADIUS =	4.700 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.02425531	WATTS/CM ²
INTERVAL POINT:	RADIUS =	4.800 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.02357193	WATTS/CM ²
INTERVAL POINT:	RADIUS =	4.900 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.022338461	WATTS/CM ²
INTERVAL POINT:	RADIUS =	5.000 CM	THETA =	0.0 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.02085281	WATTS/CM ²
INTERVAL POINT:	RADIUS =	0.100 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.016769744	WATTS/CM ²
INTERVAL POINT:	RADIUS =	0.200 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.0166665903	WATTS/CM ²
INTERVAL POINT:	RADIUS =	0.300 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.016197692	WATTS/CM ²
INTERVAL POINT:	RADIUS =	0.400 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.014301543	WATTS/CM ²
INTERVAL POINT:	RADIUS =	0.500 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.012990664	WATTS/CM ²
INTERVAL POINT:	RADIUS =	0.600 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.011532387	WATTS/CM ²
INTERVAL POINT:	RADIUS =	0.700 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.008057005	WATTS/CM ²
INTERVAL POINT:	RADIUS =	0.800 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.006490502	WATTS/CM ²
INTERVAL POINT:	RADIUS =	0.900 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.007053386	WATTS/CM ²
INTERVAL POINT:	RADIUS =	1.000 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.007651502	WATTS/CM ²
INTERVAL POINT:	RADIUS =	1.100 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.0066638423	WATTS/CM ²
INTERVAL POINT:	RADIUS =	1.200 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.003799579	WATTS/CM ²
INTERVAL POINT:	RADIUS =	1.300 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.003171812	WATTS/CM ²
INTERVAL POINT:	RADIUS =	1.400 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.00278713	WATTS/CM ²
INTERVAL POINT:	RADIUS =	1.500 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.002619987	WATTS/CM ²
INTERVAL POINT:	RADIUS =	1.600 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.002666781	WATTS/CM ²
INTERVAL POINT:	RADIUS =	1.700 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.002860828	WATTS/CM ²
INTERVAL POINT:	RADIUS =	1.800 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.003164644	WATTS/CM ²
INTERVAL POINT:	RADIUS =	1.900 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.003529164	WATTS/CM ²
INTERVAL POINT:	RADIUS =	2.000 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.003905601	WATTS/CM ²
INTERVAL POINT:	RADIUS =	2.100 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.004253939	WATTS/CM ²
INTERVAL POINT:	RADIUS =	2.200 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.004538811	WATTS/CM ²
INTERVAL POINT:	RADIUS =	2.300 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.004728578	WATTS/CM ²
INTERVAL POINT:	RADIUS =	2.400 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.004811481	WATTS/CM ²
INTERVAL POINT:	RADIUS =	2.500 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.0047818952	WATTS/CM ²
INTERVAL POINT:	RADIUS =	2.600 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.004659416	WATTS/CM ²
INTERVAL POINT:	RADIUS =	2.700 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.004446792	WATTS/CM ²
INTERVAL POINT:	RADIUS =	2.800 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.004165390	WATTS/CM ²
INTERVAL POINT:	RADIUS =	2.900 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.003842898	WATTS/CM ²
INTERVAL POINT:	RADIUS =	3.000 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.003520598	WATTS/CM ²
INTERVAL POINT:	RADIUS =	3.100 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.003210761	WATTS/CM ²
INTERVAL POINT:	RADIUS =	3.200 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.002942315	WATTS/CM ²
INTERVAL POINT:	RADIUS =	3.300 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.002733984	WATTS/CM ²
INTERVAL POINT:	RADIUS =	3.400 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.0025906017	WATTS/CM ²
INTERVAL POINT:	RADIUS =	3.500 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.00253577	WATTS/CM ²
INTERVAL POINT:	RADIUS =	3.600 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.002556403	WATTS/CM ²
INTERVAL POINT:	RADIUS =	3.700 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.002644026	WATTS/CM ²
INTERVAL POINT:	RADIUS =	3.800 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.002780315	WATTS/CM ²
INTERVAL POINT:	RADIUS =	3.900 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.002953362	WATTS/CM ²
INTERVAL POINT:	RADIUS =	4.000 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.003150973	WATTS/CM ²
INTERVAL POINT:	RADIUS =	4.100 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.003343567	WATTS/CM ²
INTERVAL POINT:	RADIUS =	4.200 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.003523249	WATTS/CM ²
INTERVAL POINT:	RADIUS =	4.300 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.003670430	WATTS/CM ²
INTERVAL POINT:	RADIUS =	4.400 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.003777670	WATTS/CM ²
INTERVAL POINT:	RADIUS =	4.500 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.003840348	WATTS/CM ²
INTERVAL POINT:	RADIUS =	4.600 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.003859740	WATTS/CM ²
INTERVAL POINT:	RADIUS =	4.700 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.003837654	WATTS/CM ²
INTERVAL POINT:	RADIUS =	4.800 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.003784656	WATTS/CM ²
INTERVAL POINT:	RADIUS =	4.900 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.003710009	WATTS/CM ²
INTERVAL POINT:	RADIUS =	5.000 CM	THETA =	180.00 DEG	PHI =	0.0 DEG	ASOSRBD	POWER	DENSITY =	0.003618500	WATTS/CM ²

APPROXIMATE EXECUTION TIME = 0.17 CPU MINUTE

The plot of the computed absorbed power densities at internal points of the sphere, Figure B-1, is similar to that in Kritikos (2,p.58) for a sphere of radius 5 cm and 1000 MHz. The electrical properties ϵ and σ are those of biological tissue of high water content.

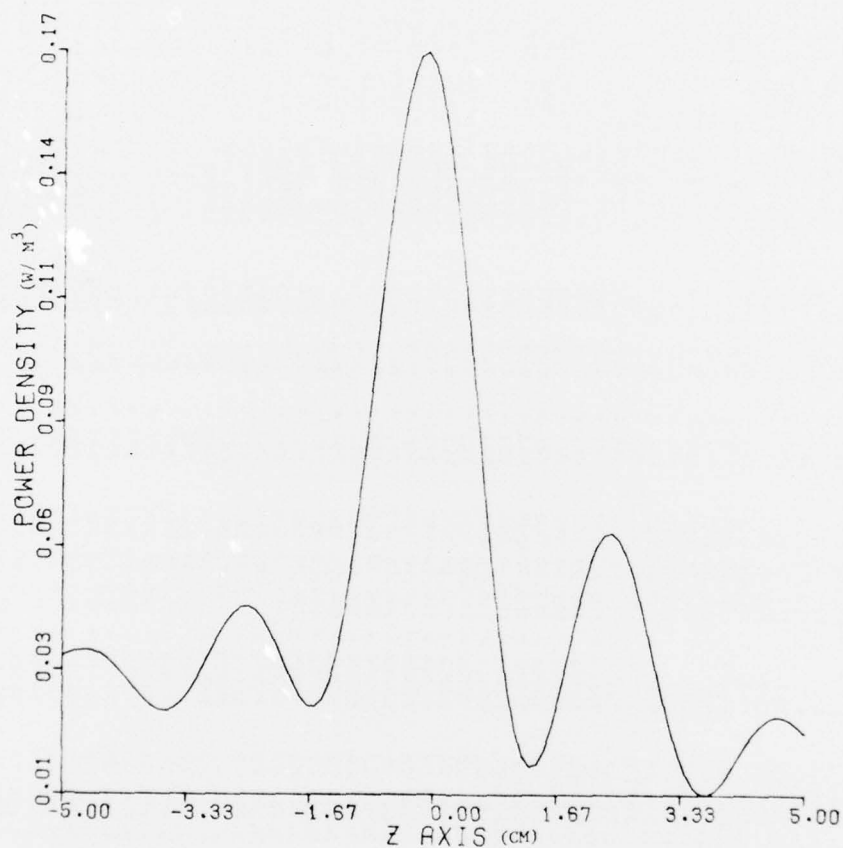


Figure B-1. Distribution of the power density inside the sphere along the z axis.

APPENDIX C

SOURCE LISTING OF PROGRAM MIE

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C                                     PROGRAM MIM                                     MIE0001
C                                     DEPOSITION OF POWER INSIDE OF A HOMOGENEOUS SPHERE MIE0002
C                                     IMMERSED IN AN ELECTROMAGNETIC FIELD MIE0003
C                                     MIE0004
C                                     MIE0005
C                                     MIE0006
C                                     MIE0007
C                                     MIE0008
C                                     MIE0009
C                                     MIE0010
C                                     MIE0011
C                                     MIE0012
C                                     MIE0013
C                                     MIE0014
C                                     MIE0015
C                                     MIE0016
C                                     MIE0017
C                                     MIE0018
C                                     MIE0019
C                                     MIE0020
C                                     MIE0021
C                                     MIE0022
C                                     MIE0023
C                                     MIE0024
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C                                     MIE0026
C                                     MIE0027
C                                     MIE0028
C                                     MIE0029
C                                     MIE0030
C                                     MIE0031
C                                     MIE0032
C                                     MIE0033
C                                     MIE0034
C                                     MIE0035
C                                     MIE0036
C                                     MIE0037
C                                     MIE0038
C                                     MIE0039
C                                     MIE0040
C                                     MIE0041
C                                     MIE0042
C                                     MIE0043
C                                     MIE0044
C                                     MIE0045
C                                     MIE0046
C                                     MIE0047
C                                     MIE0048
C                                     MIE0049
C                                     MIE0050
C                                     MIE0051
C                                     MIE0052
C                                     MIE0053
C                                     MIE0054
C                                     MIE0055
C                                     MIE0056
C                                     MIE0057
C                                     MIE0058
C                                     MIE0059
C                                     MIE0060

IMPLICIT REAL*8 (A-H,O-Z)
COMMON /GAUSS/D,M1,M2,N
COMMON CM(100),DM(100),AJR(100),AM,ANK,CEX,Z,DP(100),P(101),ALAMDA
1,DI,EO,DHI,DIH,EKE,STOPE,SLNTH,THETA,SC,NM2
COMPLEX*16 CM,DV,AJR,AM,ANK,CEX,Z,CD,Q,CPIH,CPIH,PSI,PSI
PIE=3.141592653589793D0
RAD=180.D0/PIE
STOPE = 1.D05
OLETH=999.D0
OLDD1=1.D-70
5 READ (5,10,END=80) FREQ,EPS,SIGM1,EO,D,TIME,MOC,LOFT,M1,M2,M3
10 FORMAT (6E10.0,2I5,3I3)
DOMEGA=2.D6*PIE*FREQ
SIGMA=SIGM1*9.D0
AM=CDSQRT(DCHPLX(EPS,(-4.D0*PIE*SIGMA)/DOMEGA))
ARG=DOMEGA*TIME
CEX=DCHPLX(DCOS(ARG),DSIN(ARG))
ALAMDA=2.997924562D4/FREQ
ANK=2.D0*AM*PIE/ALAMDA
EO=EO*.3333333333333333D-4
X=PIE*D/ALAMDA
Z=AM*X
C *** GENERATE SPHERICAL BESSEL FUNCTIONS JN(MX) AND NN(MX)
CALL BESS(AJR,Z,N,CQ,STOPE)
C *** GENERATE SPHERICAL BESSEL FUNCTIONS JN(X) AND NN(X)
DP(1)=-DCOS(X)/X
DP(2)=DP(1)/X-DSIN(X)/X
DO 15 M=2,N
AN2=2*M-1
DP(M+1)=AN2/X*DP(M)-DP(M-1)
IF (DABS(DP(M+1)).GE.STOPE) GO TO 20
15 CONTINUE
20 NP1=M+1
N=M
P(NP1)=0.D0
P(NP1-1)=-1.D0/(X*X*DP(NP1))
NM2=NP1-2
DO 25 M=1,NM2
I=NP1-M
AN2=2*I-1
25 P(I-1)=AN2/X*P(I)-P(I+1)
RQ=DSIN(X)/(X*P(1))
DO 30 M=1,NP1
30 P(M)=RQ*P(M)
NM1=N-1
NM2=N-2
CQ=1.D0/CQ
RQ=1.D0/RQ
C *** PRINT OUT TITLE, BASIC INPUT DATA, AND ERROR ANALYSIS DATA
PRINT 35,FREQ,ALAMDA,EO,SIGM1,EPS,D,TIME,AM,NM1,X,RQ,CQ
35 FORMAT ('DEPOSITION OF POWER INSIDE A HOMOGENEOUS SPHERE IMMERSED IN AN ELECTROMAGNETIC FIELD',F10.2,' MHZ WAVELENGTH',F15.5,' CM FIELD STRENGTH',F7.2,' V/M',F10.6,' MEO/M RELATIVE DIELECTRIC CONSTANT',F7.2,' MIE0058
1 DIAMETER',F7.2,' CM TIME',F7.2,' SEC',F7.2,' MIE0059
1 INDEX',F12.5,' OTHER HIGHEST ORDER NEUMANN FUNCTION COMPUTED IS',F7.2,' MIE0060

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1 OF ORDER =',I3,' SIZE PARAMETER =',E13.5/'CRATIO OF THE FIRSTNIE0061
1 CALCULATION OF THE ZERO ORDER SPHERICAL BESSEL FUNCTION OF ARGUMENTNIE0062
1T X/' TO THE VALUE OF SIN(X)/X =',E15.5/'CRATIO OF THE FIRST CALCNIE0063
1OLATION OF THE ZERO ORDER SPHERICAL BESSEL FUNCTION OF ARGUMENT Z/'NIE0064
1' TO THE VALUE OF SIN(Z)/Z =',2E15.5/'NIE0065
C *** GENERATE COEFFICIENTS CN AND DN NIE0066
DO 40 NN=1,NM2 NIE0067
NP1=NN+1 NIE0068
RPHI=X*P(NP1) NIE0069
RPSI=DCMLX(RPHI,-X*DP(NP1)) NIE0070
CPHI=Z*AJR(NP1) NIE0071
AN=NN NIE0072
FACT=1.D0/(2.D0*AN+1.D0) NIE0073
RPHIP=P(NP1)+X*(FACT*(AN*P(NP1-1)-(AN+1.D0)*P(NP1+1))) NIE0074
RKIP=-DP(NP1)-X*(FACT*(AN*DP(NP1-1)-(AN+1.D0)*DP(NP1+1))) NIE0075
FPSIP=DCMLX(RPHIP,RKIP) NIE0076
CPHIP=AJR(NP1)+Z*(FACT*(AN*AJR(NP1-1)-(AN+1.D0)*AJR(NP1+1))) NIE0077
Q=(RPHIP*RPSI-RPHI*RPSIP) NIE0078
CN(NN)=Q/(CPHIP*RPSI-AN*CPHI*RPSIP) NIE0079
40 DN(NN)=Q/(AN*CPHIP*RPSI-CPHI*RPSIP) NIE0080
IF (NOC.EQ.0) GO TO 65 NIE0081
DO 60 J=1,NOC NIE0082
READ 10, R,THETAD,PHID NIE0083
D1=2.D0*R NIE0084
PHI=PHID/RAD NIE0085
IF (THETAD.EQ.OLDTH) GO TO 45 NIE0086
THETA=THETAD/RAD NIE0087
SINTH=DSIN(THETA) NIE0088
CALL PL(THETA,N,P,DP) NIE0089
45 IF (D1.EQ.OLDD1) GO TO 50 NIE0090
RKR=ALAMDA/(PIE*D1) NIE0091
Z=AN*PIE*D1/ALAMDA NIE0092
CALL BESS(AJR,Z,NC,Q,STOPR) NIE0093
NC=NC-2 NIE0094
IF (NC.GT.NM2)NC=NM2 NIE0095
C ABSORBED POWER DENSITY AT GIVEN INTERNAL POINT OF SPHERE NIE0096
50 CALL EVEC(PD) NIE0097
PD=.05D1*SIGMA*PD NIE0098
C *** PRINT OUT INTERIOR POINT PARTICULARS NIE0099
PRINT 55,R,THETAD,PHID,PD NIE0100
55 FORMAT (' INTERNAL POINT: RADIUS =',F8.3,' CM .THETA =',F7.2,' NIE0101
1DEG PHI =',F7.2,' DEG ABSORBED POWER DENSITY =',F12.8,' WATNIE0102
1TS/M**3') NIE0103
60 CONTINUE NIE0104
65 IF (IOPT.EQ.0) GO TO 5 NIE0105
C TOTAL ABSORBED POWER AND AVERAGE ABSORBED POWER DENSITY NIE0106
TOTPOW=.05D-6*SIGMA*GAUSS3(M3) NIE0107
PAVG=TOTPOW*1.D6/((4.D0/3.D0)*PIE*(D/2.D0)**3) NIE0108
C *** PRINT OUT ABSORBED TOTAL POWER AND AVERAGE ABSORBED NIE0109
POWER DENSITY NIE0110
PRINT 75,TOTPOW,PAVG NIE0111
75 FORMAT ('0',9X,'TOTAL ABSORBED POWER=',1PE13.5,' WATTS'/'0',9X,'AVNIE0112
1ERAGE ABSORBED POWER DENSITY=',E13.5,' WATTS/M**3') NIE0113
GO TO 5 NIE0114
80 STOP NIE0115
END NIE0116
SUBROUTINE EVEC(PD) NIE0117
C COMPUTES RADIAL,COLATITUDE,AND AZIMUTHAL COMPONENTS OF NIE0118
ELECTRIC FIELD VECTOR E AND THE SCALAR PRODUCT E.E* NIE0119
C IMPLICIT REAL*8 (A-H,O-Z) NIE0120

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COMMON CN(100),DN(100),AJR(100),AM,ANK,CEX,Z,DP(100),P(101),ALAMPAMIE0121
1,DI,EO,PHI,PIE,BKR,STOPE,SINTE,THETA,NC,NN2 MIE0122
COMPLEX*16 CN,DN,AJR,AM,ANK,CEX,Z,UE,UT,UP,URE,URT,URP,URRE,URRT,UMIE0123
1RRP,DERJN1,AMZ,SUM,ER,ETHETA,EPHI,TE,TT,TP,T1,T2,T3,ZSQ MIE0124
DATA EPS/1.D-10/ MIE0125
ZSQ=Z**Z MIE0126
UR=DCMPLX(0.D0,0.D0) MIE0127
URE=DCMPLX(0.D0,0.D0) MIE0128
URRE=DCMPLX(0.D0,0.D0) MIE0129
UT=DCMPLX(0.D0,0.D0) MIE0130
URT=DCMPLX(0.D0,0.D0) MIE0131
URRT=DCMPLX(0.D0,0.D0) MIE0132
UP=DCMPLX(0.D0,0.D0) MIE0133
URP=DCMPLX(0.D0,0.D0) MIE0134
URRP=DCMPLX(0.D0,0.D0) MIE0135
IE=0 MIE0136
IT=0 MIE0137
IP=0 MIE0138
NCK=1 MIE0139
DO 70 NN=1,NC MIE0140
NNP1=NN+1 MIE0141
FAC1 = 2*NN+1 MIE0142
FAC2=NN*FAC1 MIE0143
FAC3=FAC1/FAC2 MIE0144
DERJN1=(NN*AJR(NN)-NNP1*AJR(NN+2))/FAC1 MIE0145
IF (IE.EQ.1) GO TO 5 MIE0146
TE=CN(NN)*FAC3*P(NN) MIE0147
T1=TE*AJR(NN+1) MIE0148
CALL TERM(NCK,T1) MIE0149
UE=UE+T1 MIE0150
T2=TE*DERJN1 MIE0151
CALL TERM(NCK,T2) MIE0152
URE=URE+T2 MIE0153
T3=TE*(-2.D0*DERJN1/Z-(1.D0-FAC2/ZSQ)*AJR(NN+1)) MIE0154
CALL TERM(NCK,T3) MIE0155
URRE=URRE+T3 MIE0156
IF (CDABS(T1).GT.CDABS(UE)*EPS) GO TO 5 MIE0157
IF (CDABS(T2).GT.CDABS(URE)*EPS) GO TO 5 MIE0158
IF (CDABS(T3).LE.CDABS(URRE)*EPS) IE=1 MIE0159
5 IF (IT.EQ.1) GO TO 45 MIE0160
IF ((THETA.GE.1.0D-6).AND.(THETA.LT.3.141592D0)) GO TO 20 MIE0161
TT=DN(NN)*FAC1/2.D0 MIE0162
IF (THETA.GE.3.141592D0) TT=(-1.D0)**NNP1*TT MIE0163
GO TO 25 MIE0164
20 TT=DN(NN)*FAC3*P(NN)/SINTE MIE0165
25 T1=TT*AJR(NN+1) MIE0166
CALL TERM(NCK,T1) MIE0167
UT=UT+T1 MIE0168
TT=CN(NN)*FAC3*DP(NN) MIE0169
T2=TT*AJR(NN+1) MIE0170
CALL TERM(NCK,T2) MIE0171
URT=URT+T2 MIE0172
T3=TT*DERJN1 MIE0173
CALL TERM(NCK,T3) MIE0174
URRT=URRT+T3 MIE0175
IF (CDABS(T1).GT.CDABS(UT)*EPS) GO TO 45 MIE0176
IF (CDABS(T2).GT.CDABS(URT)*EPS) GO TO 45 MIE0177
IF (CDABS(T3).LE.CDABS(URRT)*EPS) IT=1 MIE0178
45 IF (IP.EQ.1) GO TO 68 MIE0179
TP=DN(NN)*FAC3*DP(NN) MIE0180

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T1=TP*AJB(NN+1)
CALL TERM(NCK,T1)
UP=UP+T1
IF ((THETA.GE.1.0D-6).AND.(THETA.LT.3.141592D0)) GO TO 60
TP=CN(NN)*PAC1/2.D0
IF (THETA.GE.3.141592D0) TP=(-1.D0)**NNPI*TP
GO TO 65
60 TP=CN(NN)*PAC3*P(NN)/SINTH
65 T2=TP*AJB(NN+1)
CALL TERM(NCK,T2)
URP=URP+T2
T3=TP*DERON1
CALL TERM(NCK,T3)
URRP=URRP+T3
IF (CDABS(T1).GT.CDABS(UP)*EPS) GO TO 68
IF (CDABS(T2).GT.CDABS(URP)*EPS) GO TO 68
IF (CDABS(T3).LE.CDABS(URRP)*EPS) IP=1
68 IF (I2+T2+IP.2Q.3) GO TO 71
NCK=NCK+1
71 IF (NCK.GT.4) NCK=1
71 PCO=DCOS(PHI)
AMZ=AM*Z
ER=CEX*PCO*(2.D0*AM*URB+AMZ*(URR+UR))
ER=EO*DCMPLX(-DIMAG(ER),DREAL(ER))
SUM=CEX*PCO*(RKR*URT+AM*URRT)
ETHETA=EO*(AM*CEX*PCO*UT+DCMPLX(-DIMAG(SUM),DREAL(SUM)))
PSI=DSIN(PHI)
SUM=-CEX*PSI*(RKR*URP+AM*URRP)
EPHI=EO*(-AM*CEX*PSI*UP+DCMPLX(-DIMAG(SUM),DREAL(SUM)))
PD=DREAL(ER*DCONJG(ER))+DREAL(ETHETA*DCONJG(ETHETA))+DREAL(EPHI*DCONJG(EPHI))
10NUG(EPHI))
RETURN
END
SUBROUTINE TERM(NCK,T)
C      COMPUTES (-J)**N*(N-TH TERM IN SERIES)
IMPLICIT REAL*8 (A-H,O-Z)
COMPLEX*16 T
GO TO (5,10,15,20),NCK
5 T=DCMPLX(DIMAG(T),-DREAL(T))
GO TO 20
10 T=-T
GO TO 20
15 T=DCMPLX(-DIMAG(T),DREAL(T))
20 RETURN
END
SUBROUTINE BESS(AJ,Z,N,Q,STOPR)
C      GENERATES SPHERICAL BESSEL FUNCTIONS JN(MX) AND NN(MX)
IMPLICIT REAL*8 (A-H,O-Z)
DIMENSION AJ(100)
COMPLEX*16 Z,AJ,Y,Y0,Y1,Q
Y0 =-CDCOS(Z)/Z
Q=CDSIN(Z)
Y1 = (Y0-Q)/Z
DO 5 M=3,100
Y=(2*M-3)/Z*Y1-Y0
IF (CDABS(Y).GE.STOPR) GO TO 10
Y0=Y1
5 Y1=Y
10 IF (M.GT.3) GO TO 30
C *** PRINT OUT ERROR MESSAGE

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      PRINT 25,Z
25  FORMAT ('0**** PROCESS CAN NOT PROCEED SINCE NM2 = 0 FOR Z =',1P2D
115.7)
      STOP
30  AJ(M)=DCMPLX(0.DC,0.D0)
      AJ(M-1)=-1.D0/(Z*Z*Y)
      NM2=M-2
      DO 35 I=1,NM2
      L=M-I
35  AJ(L-1)=(2*L-1)/Z*AJ(L)-AJ(L+1)
      Q=Q/(Z*AJ(1))
      NM1=M-1
      DO 40 N=1,NM1
      AJ(N)=Q*AJ(N)
      IF (CDABS(AJ(N)).LT.1.D-25) RETURN
40  CONTINUE
      RETURN
      END
      SUBROUTINE PL(THETA,M,P,DP)
C      GENERATES ASSOCIATED LEGENDRE FUNCTIONS AND THEIR
C      FIRST DERIVATIVES
      IMPLICIT REAL*8 (A-H,O-Z)
      DIMENSION P(100),DP(100)
      SNJ=DSIN(THETA)
      CNJ=DCOS(THETA)
      P(1)=SNJ
      P(2)=3.D0*SNJ*CNJ
      DP(1)=CNJ
      DO 10 N=2,M
      A=M
      MP1=M+1
      P(M+1)=(2.D0*A+1.D0)/A*CNJ*P(M)-(A+1.D0)/A*P(M-1)
      IF ((THETA.GE.1.0D-6).AND.(THETA.LT.3.141592D0)) GO TO 5
      DP(M)=M*MP1/2
      IF (THETA.GE.3.141592D0) DP(M)=(-1.D0)**M*DP(M)
      GO TO 10
5    DP(M)=(A*CNJ*P(M)-(A+1.D0)*P(M-1))/SNJ
10   CONTINUE
      RETURN
      END
      FUNCTION GAUSS3(M3)
C      PRODUCT RULE: M-POINT GAUSS-LEGENDRE QUADRATURE FORMULAS
      IMPLICIT REAL*8 (A-H,O-Z)
      COMMON /GAUSS/D,M1,M2,N
      COMMON CN(100),DN(100),AJR(100),AM,ANK,CEX,Z,DP(100),P(101),ALANDAM
1,D1,ZO,PHI,PIE,RKR,STORP,SINTH,THETA,NC,NM2
      COMPLEX*16 CN,DN,AJR,AM,ANK,CEX,Z,O
      DIMENSION NPOINT(9),KEY(10),Y(33),XT(33),ARG2(2),ARG3(2)
      DATA NPOINT/2,3,4,5,6,8,10,12,14/
      DATA KEY/1,2,4,6,9,12,16,21,27,34/
      DATA Y/0.57735026918962D0,0.00000000000000D0,0.77459666924148D0,0.
133998104358486D0,0.86113631159465D0,0.00000000000000D0,0.538469310
110568D0,0.90617984593866D0,0.23861918608320D0,0.66120938646626D0,0.
1.93246951420315D0,0.18343464249565D0,0.52553240991633D0,0.79666647
1741363D0,0.96028985649754D0,0.14887433898163D0,0.43339539412925D0,
10.67940956829902D0,0.86506336668898D0,0.97390652851717D0,0.1252334
10851147D0,0.36783149899818D0,0.58731795428662D0,0.76990267419430D0,
1,0.90411725637047D0,0.98156063424672D0,0.10805494870734D0,0.319112
136892789D0,0.51524863635815D0,0.68729290481169D0,0.82720131506976D0,
10,0.92843488366357D0,0.98628380869681D0/
      M10300

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DATA WT/1.0000000000000000,0.8888888888888888,0.5555555555555555,0.0000000000000000,MIE0301
1.65214515486255D0,0.34785484513745D0,0.5688888888888888,0.478628678120302
1049937D0,0.23692688505619D0,0.46791393457269D0,0.36076157304814D0,MIE0303
10.17132449237917D0,0.36268378337836D0,0.31370664587789D0,0.2223810MIE0304
13445337D0,0.10122853629038D0,0.29552422471475D0,0.26926671931000D0MIE0305
1,0.21908636251598D0,0.14945134915058D0,0.06667134430869D0,0.249147MIE0306
104581340D0,0.23349253653835D0,0.20316742672307D0,0.16007832854335DMIE0307
10,0.10693932599532D0,0.04717533638651D0,0.21526385346316D0,0.20519MIE0308
1846372130D0,0.18553839747794D0,0.15720316715819D0,0.1215185706879CMIE0309
100,0.08015878715976D0,0.03511946033175D0/MIE0310
DO 5 I=1,9MIE0311
IF (M3.EQ.NPOINT(I)) GO TO 20MIE0312
5 CONTINUEMIE0313
10 PRINT 15,M1,M2,M3MIE0314
15 FORMAT ('-ERROR IN INTEGRATION CONTROLS. M1 =',I6,' M2 =',I6,'MIE0315
1M3 =',I6)MIE0316
GAUSS3=0.D0MIE0317
RETURNMIE0318
20 JF3 =KEY(I)MIE0319
JL3 =KEY(I+1)-1MIE0320
DO 25 I=1,9MIE0321
IF (M2.EQ.NPOINT(I)) GO TO 30MIE0322
25 CONTINUEMIE0323
GO TO 10MIE0324
30 JF2 =KEY(I)MIE0325
JL2 =KEY(I+1)-1MIE0326
DO 35 I=1,9MIE0327
IF (M1.EQ.NPOINT(I)) GO TO 40MIE0328
35 CONTINUEMIE0329
GO TO 10MIE0330
40 JF1 =KEY(I)MIE0331
JL1 =KEY(I+1)-1MIE0332
C INTEGRATE OVER THETAMIE0333
PD2=PI2/2.D0MIE0334
R=D/4.D0MIE0335
SUM3=0.D0MIE0336
DO 85 J3=JF3,JL3MIE0337
IF (Y(J3).NE.0.D0) GO TO 45MIE0338
ARG3(1)=PD2MIE0339
I3=1MIE0340
GO TO 50MIE0341
45 ARG3(1)=PD2+PD2*Y(J3)MIE0342
ARG3(2)=PD2-PD2*Y(J3)MIE0343
I3=2MIE0344
50 DO 85 L=1,I3MIE0345
THETA=ARG3(L)MIE0346
SINCH=DSIN(THETA)MIE0347
CALL PL(THETA,N,P,DP)MIE0348
C INTEGRATE OVER RADIUSMIE0349
SUM2=0.D0MIE0350
DO 80 J2=JF2,JL2MIE0351
IF (Y(J2).NE.0.D0) GO TO 55MIE0352
ARG2(1)=RMIE0353
I2=1MIE0354
GO TO 60MIE0355
55 ARG2(1)=R+Y(J2)*RMIE0356
ARG2(2)=R-Y(J2)*RMIE0357
I2=2MIE0358
60 DO 80 I=1,I2MIE0359
D1=2.D0*ARG2(I)MIE0360

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	RKP=ALAMDA/(PIE*D1)		
	Z=AM*PIE*D1/ALAMDA		
	CALL BESS (AJR,Z,NC,Q,STOPR)		
	NC=NC-2		
	IF (NC.GT.NM2) NC=NM2		
C	INTEGRATE OVER PHI		
	SUM1=0.D0		
	DO 75 J=JF1,JL1		
	IF (Y(J).NE.0.D0) GO TO 65		
	PHI=PIE		
	XMUL1=1.D0		
	GO TO 70		
65	PHI=PIE*Y(J)*PIE		
C	FUNC(180+PHI)=FUNC(180-PHI)		
	XMUL1=2.D0		
70	CALL EVEC(PD)		
75	SUM1=SUM1+XMUL1*WT(J)*PD		
80	SUM2=SUM2+WT(J2)*SUM1*ARG2(I)**2		
85	SUM3=SUM3+WT(J3)*SUM2*SINTN		
	GAUSS3=SUM3*PD2*B*PIE		
	RETURN		
	END		

MIE0361
 MIE0362
 MIE0363
 MIE0364
 MIE0365
 MIE0366
 MIE0367
 MIE0368
 MIE0369
 MIE0370
 MIE0371
 MIE0372
 MIE0373
 MIE0374
 MIE0375
 MIE0376
 MIE0377
 MIE0378
 MIE0379
 MIE0380
 MIE0381
 MIE0382